

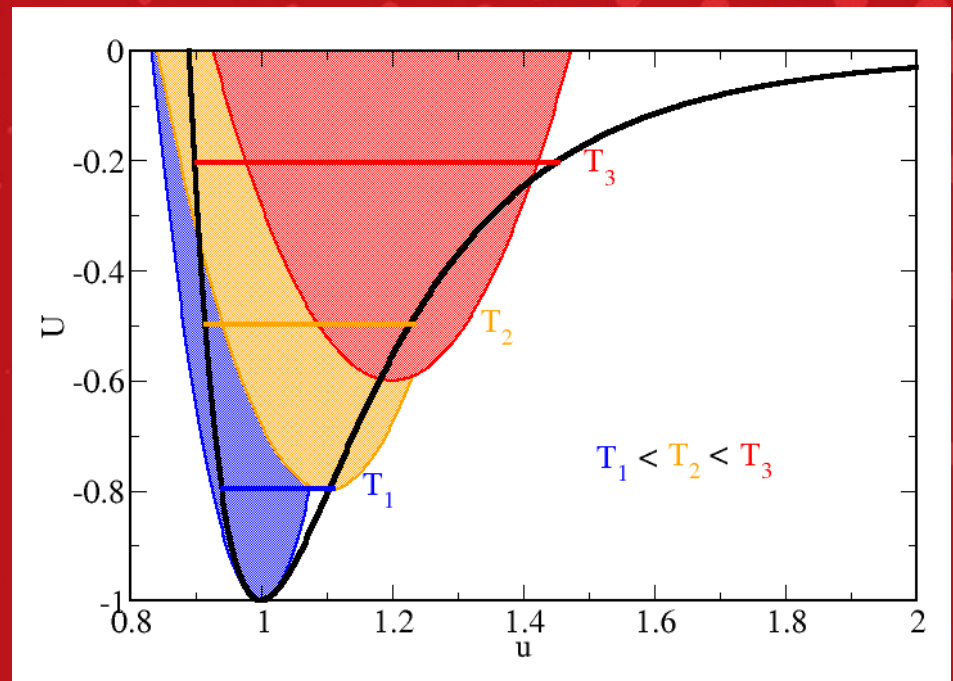
aTDEP -- Lattice dynamics including anharmonicity

F. Bottin & J. Bouchet

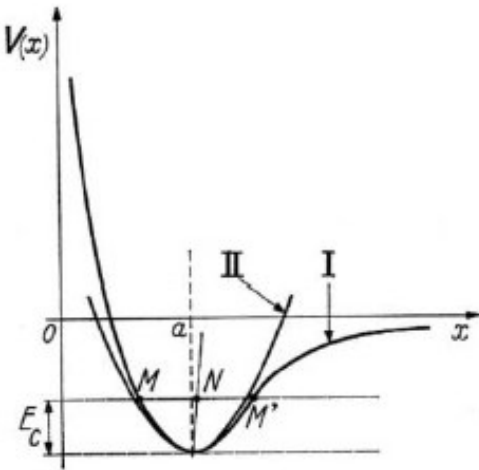
CEA, DAM, DIF, F-91297 Arpajon, France



DE LA RECHERCHE À L'INDUSTRIE



ABINIT Developer Workshop May 31 - June 4, 2021



Harmonic approximation

The temperature acts only through the filling of the energy levels.

Quasi-harmonic approximation

Includes the thermal expansion. The phonon frequencies **implicitly** depend on the temperature (through the volume).

Calculations at 0 K (DFPT, FD...)

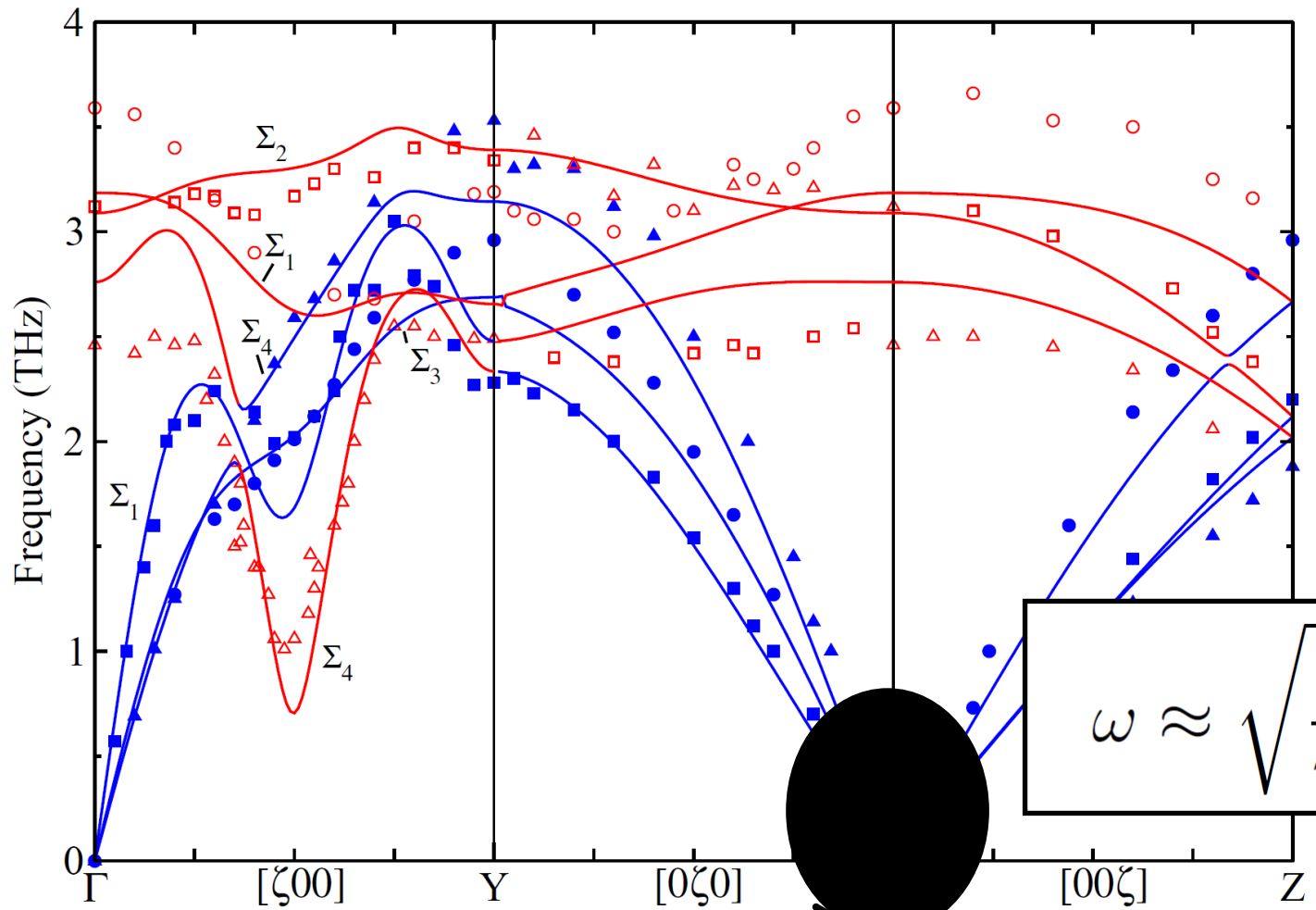
The QHA gives good results in numerous cases... except when the phonon frequencies evolve at constant V .

$$\omega(T, V(T))$$

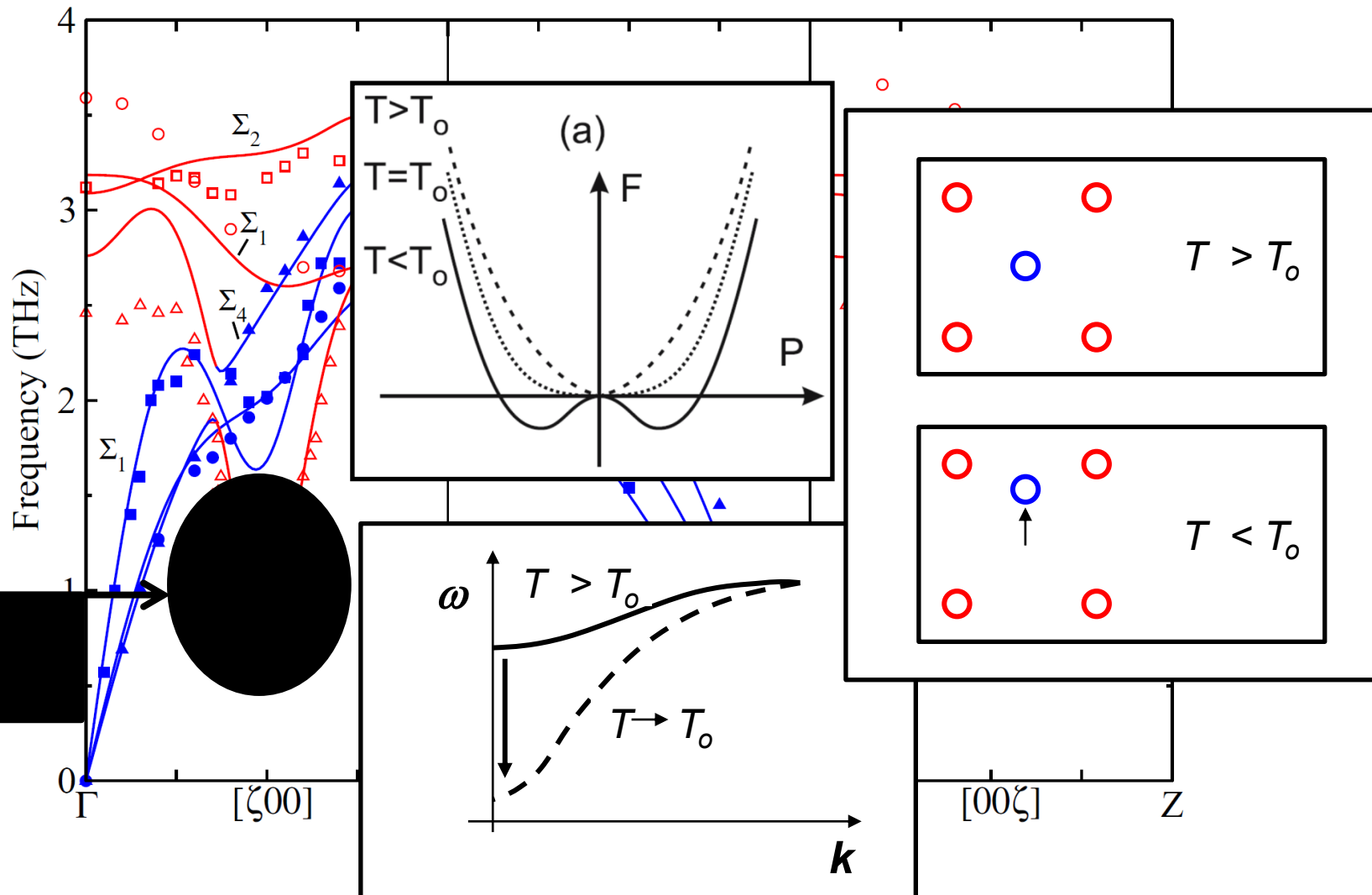
Anharmonic effects

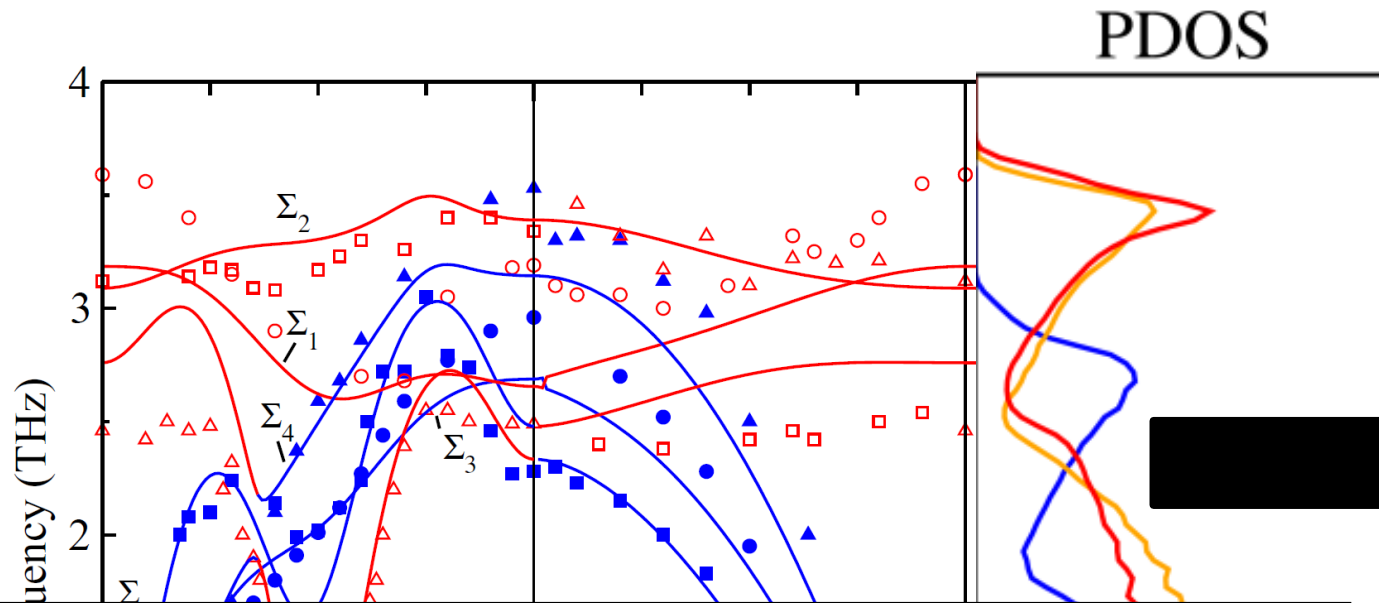
The phonon frequencies **explicitly** depend on the temperature

The temperature has to be explicitly taken into account (MD, MC...)



$$\omega \approx \sqrt{\frac{\lambda}{m}} ak$$



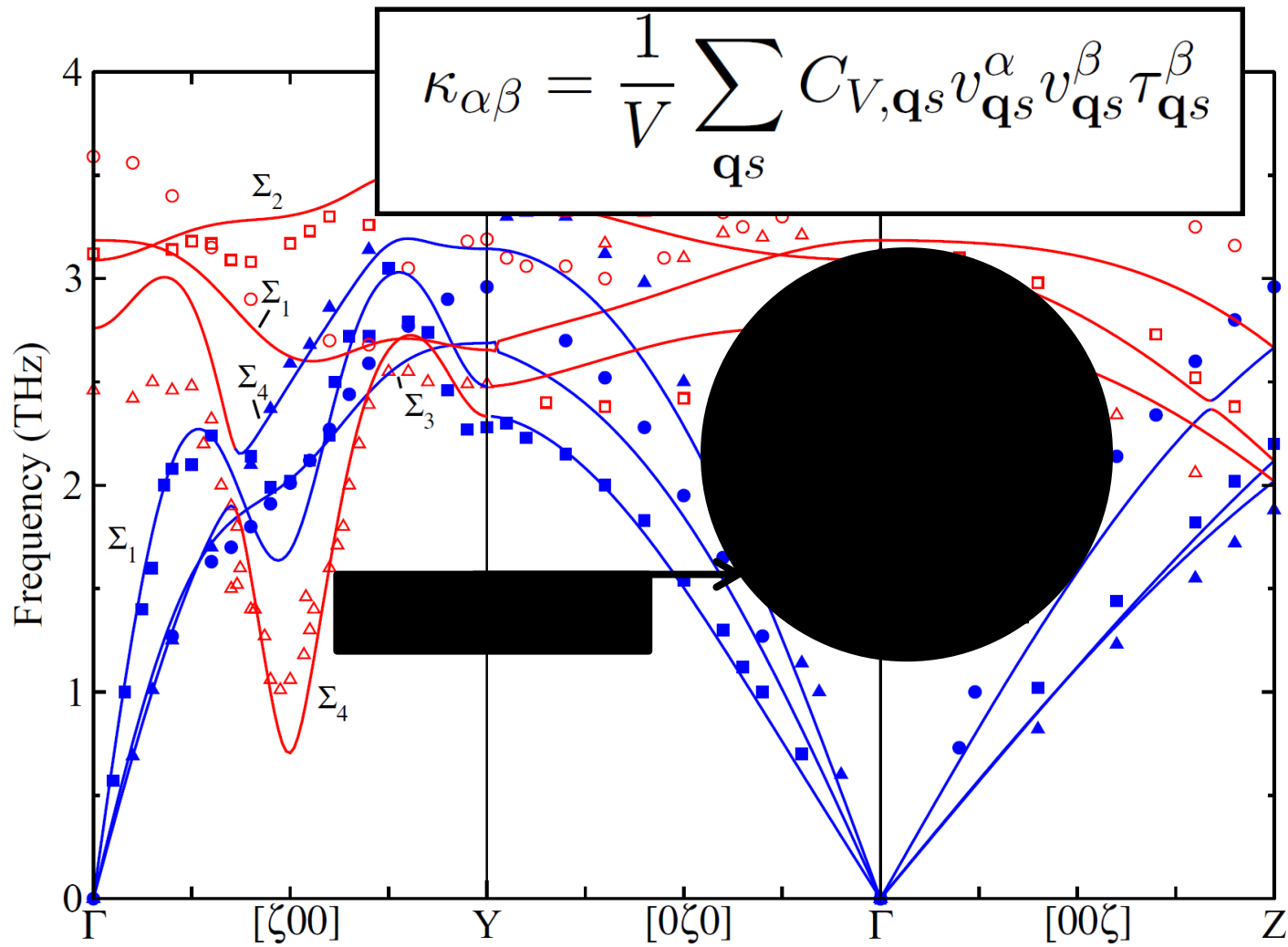


$$U = 3N_a \int_0^{\omega_{max}} \frac{\hbar\omega}{2} \coth\left(\frac{\beta\hbar\omega}{2}\right) g(\omega) d\omega$$

$$C_{vib,V} = 3N_a k_B \int_0^{\omega_{max}} \left(\frac{\beta\hbar\omega}{2 \sinh(\frac{\beta\hbar\omega}{2})} \right)^2 g(\omega) d\omega$$

$$F_{vib} = \frac{3N_a}{\beta} \int_0^{\omega_{max}} \ln \left(2 \sinh\left(\frac{\beta\hbar\omega}{2}\right) \right) g(\omega) d\omega \longrightarrow F(V, T) = U_0(V) + F_{vib}(V, T)$$

$$S_{vib} = 3N_a k_B \int_0^{\omega_{max}} \left[\frac{\beta\hbar\omega}{2} \coth\left(\frac{\beta\hbar\omega}{2}\right) - \ln \left(2 \sinh\left(\frac{\beta\hbar\omega}{2}\right) \right) \right] g(\omega) d\omega$$



Everything depends on IFC

Grüneisen parameter :

$$\gamma_i = - \left(\frac{\partial \ln \omega_i}{\partial \ln V} \right)_T = - \frac{V}{\omega_i} \left(\frac{\partial \omega_i}{\partial V} \right)_T$$

$$\gamma_s(\mathbf{q}) = - \frac{1}{6\omega_s^2(\mathbf{q})} \sum_{ijk,bc,\alpha\beta\gamma} \Psi_{ijk}^{\alpha\beta\gamma}(0,b,c) \frac{X_{is}^{\star\alpha}(\mathbf{q}) X_{js}^{\beta}(\mathbf{q})}{\sqrt{M_i M_j}} \tau_k^{\gamma} \exp[i\mathbf{q} \cdot \mathbf{R}(b)]$$

specific heat :

$$\gamma = \frac{\sum_{i=1}^{3N_a} \gamma_i C_{V,i}}{C_V} \quad \max \left(\frac{\beta \hbar \omega}{2 \sinh(\frac{\beta \hbar \omega}{2})} \right)^2 g(\omega) d\omega$$

$$g(\omega) = \frac{1}{3N_a} \sum_{i=1}^{N_a} \delta(\omega - \omega_i)$$

$$\alpha_p = \frac{\gamma C_V}{B_T V}$$

Bulk Modulus :

Thermal expansion :

$$C_{\alpha\beta\gamma\delta} = A_{\alpha\gamma\beta\delta} + A_{\beta\gamma\alpha\delta} - A_{\alpha\beta\gamma\delta}$$

$$\alpha_p = \frac{1}{B} \sum_{i=1}^{3N_a} \left(- \frac{C_{V,i}}{\omega_i} \right) \left(\frac{\partial \omega_i}{\partial V} \right)_T \frac{1}{V} \sum_{ij} \Phi_{ij}^{\alpha\beta} d_{ij}^{\gamma} d_{ij}^{\delta}$$

$$B_T = ((C_{11} + C_{22} + C_{33}) + 2(C_{12} + C_{13} + C_{23}))/9$$

Taylor expansion of the potential energy around equilibrium :

$$U_{\text{model}} = U_0 + \sum_{i,\alpha} \Pi_i^\alpha u_i^\alpha + \frac{1}{2!} \sum_{ij,\alpha\beta} \Phi_{ij}^{\alpha\beta} u_i^\alpha u_j^\beta + \frac{1}{3!} \sum_{ijk,\alpha\beta\gamma} \Psi_{ijk}^{\alpha\beta\gamma} u_i^\alpha u_j^\beta u_k^\gamma + O(u^4)$$

Here, the forces are :

$$\begin{aligned} \mathcal{F}_{i,\text{model}}^\alpha &= -\Pi_i^\alpha - \sum_{j,\beta} \Phi_{ij}^{\alpha\beta} u_j^\beta - \frac{1}{2} \sum_{jk,\beta\gamma} \Psi_{ijk}^{\alpha\beta\gamma} u_j^\beta u_k^\gamma + O(u^3) \\ &= \sum_{p\lambda} f_{i,\lambda p}^\alpha(\mathbf{u}) \theta^{\lambda p} \end{aligned}$$

Non-linear wrt u

Linear wrt θ

$$\text{AIMD : } (\mathbf{F}_{MD}(\mathbf{t}) ; \mathbf{u}_{MD}(\mathbf{t})) \longrightarrow \mathcal{F}_{i,\text{MD}}^\alpha(t) = \sum_{p\lambda} f_{i,\lambda p}^\alpha(\mathbf{u}_{MD}(t)) \theta^{\lambda p}$$

$\longrightarrow$ solved using a least squares method

The solution giving the lowest residual is :

$$\Theta = \mathbf{f}^\dagger \cdot \mathbf{F}_{MD}$$

O. Hellman *et al.*, PRB **84**, 180301(R) (2011) ; O. Hellman *et al.*, PRB **87**, 104111 (2013).

J. Bouchet & F. Bottin., Phys. Rev. B **92**, 174108 (2015) ; F. Bottin, J. Bieder & J. Bouchet CPC 254, 107301 (2020).

Using invariances (translation & rotation), symmetries of the crystal, dynamical matrix, elastic tensor, etc, the number of IFC coefficients can be strongly reduced.

For 100 atoms : → at the 2nd order : from ~ 90 000 to ~ 20

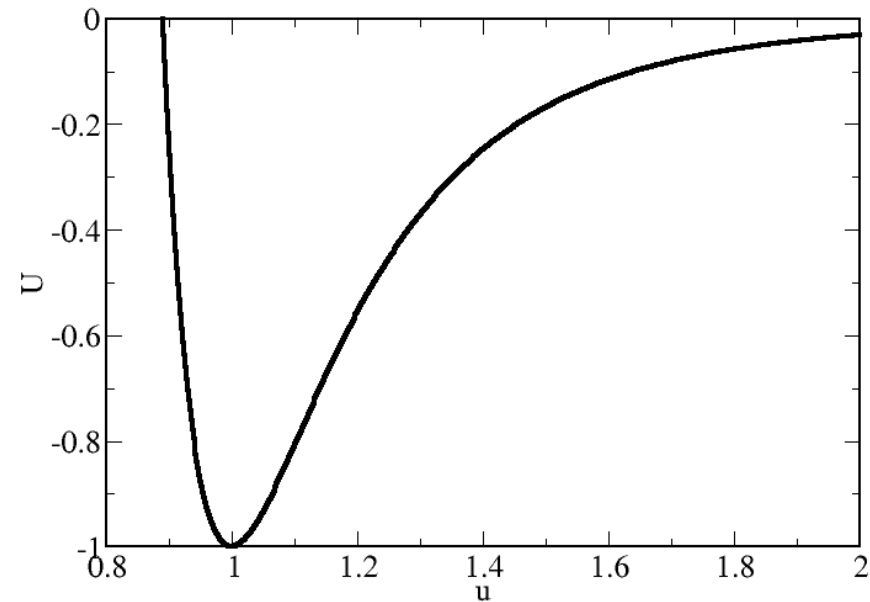
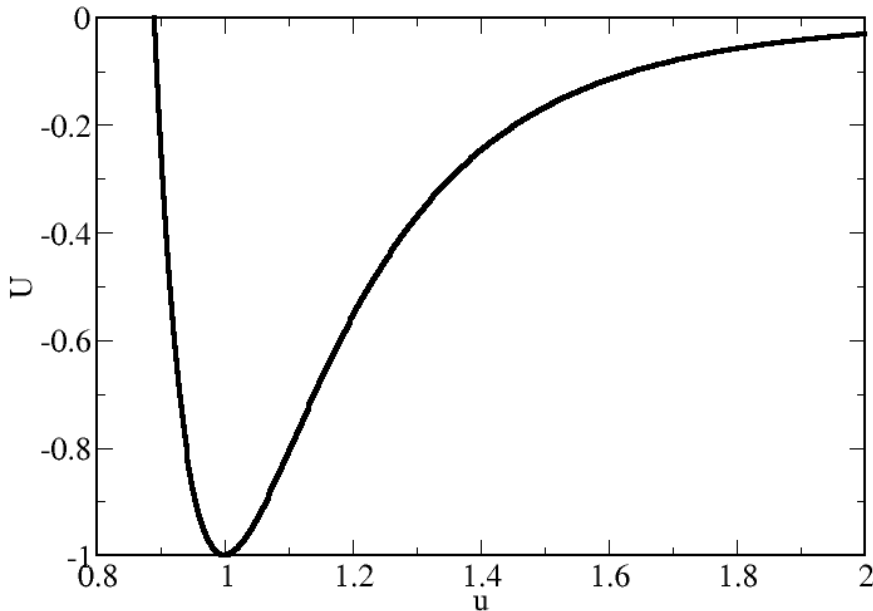
→ at the 3rd order : from ~ 27 000 000 to ~ 200

$$\begin{pmatrix} \mathcal{F}_1^x(1) \\ \vdots \\ \theta_{p\Lambda_p} \\ \vdots \\ \mathcal{F}_{N_a}^z(N_t) \end{pmatrix} = \begin{pmatrix} f_{1,11}^x(1) & f_{1,12}^x(1) & \dots & f_{1,p\Lambda_p}^x(1) \\ \vdots & \vdots & \ddots & \vdots \\ f_{N_a,11}^z(N_t) & \dots & f_{N_a,p\Lambda_p}^z(N_t) \\ \vdots & \vdots & \vdots \\ f_{N_a,11}^z(N_t) & f_{N_a,12}^z(N_t) & \dots & f_{N_a,p\Lambda_p}^z(N_t) \end{pmatrix}^\dagger \cdot \begin{pmatrix} \mathcal{F}_1^x(1) \\ \vdots \\ \mathcal{F}_{N_a}^z(N_t) \\ \vdots \\ \theta_{p\Lambda_p} \end{pmatrix}$$

$$U = U_0 + \frac{1}{2!} \sum_{ij, \alpha\beta} \Phi_{ij}^{\alpha\beta} u_i^\alpha u_j^\beta + \frac{1}{3!} \sum_{ijk, \alpha\beta\gamma} \Psi_{ijk}^{\alpha\beta\gamma} u_i^\alpha u_j^\beta u_k^\gamma + \frac{1}{4!} \sum_{ijkl, \alpha\beta\gamma\delta} \chi_{ijkl}^{\alpha\beta\gamma\delta} u_i^\alpha u_j^\beta u_k^\gamma u_l^\delta + O(u^5)$$

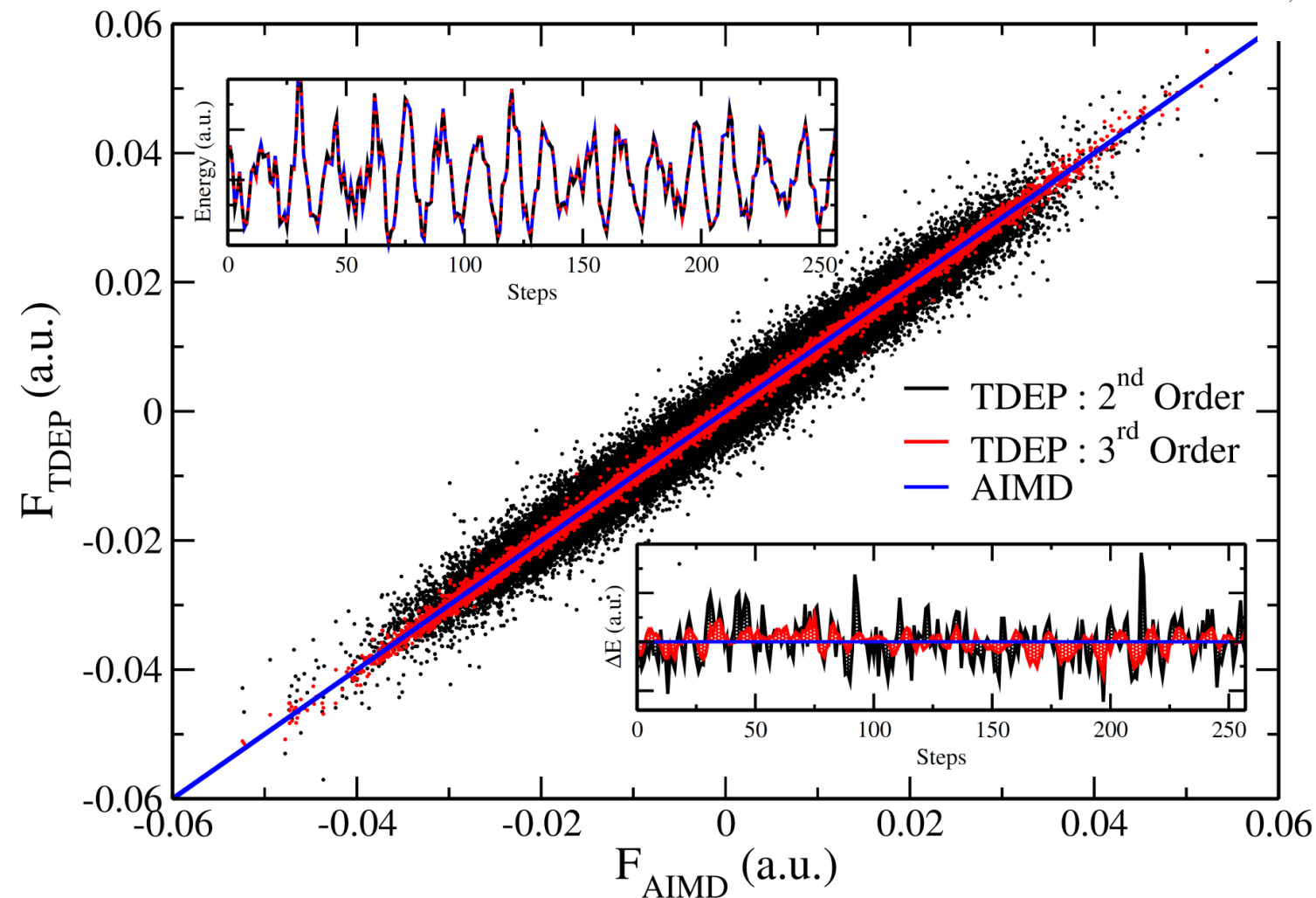
$$U = U_0 + \frac{1}{2!} \sum_{ij, \alpha\beta} \Theta_{ij}^{\alpha\beta}(T) u_i^\alpha u_j^\beta$$

$$U = U_0 + \frac{1}{2!} \sum_{ij, \alpha\beta} \Theta_{ij}^{\alpha\beta}(T) u_i^\alpha u_j^\beta + \frac{1}{3!} \sum_{ijk, \alpha\beta\gamma} \Theta_{ijk}^{\alpha\beta\gamma}(T) u_i^\alpha u_j^\beta u_k^\gamma + \frac{1}{4!} \sum_{ijkl, \alpha\beta\gamma\delta} \Theta_{ijkl}^{\alpha\beta\gamma\delta}(T) u_i^\alpha u_j^\beta u_k^\gamma u_l^\delta$$

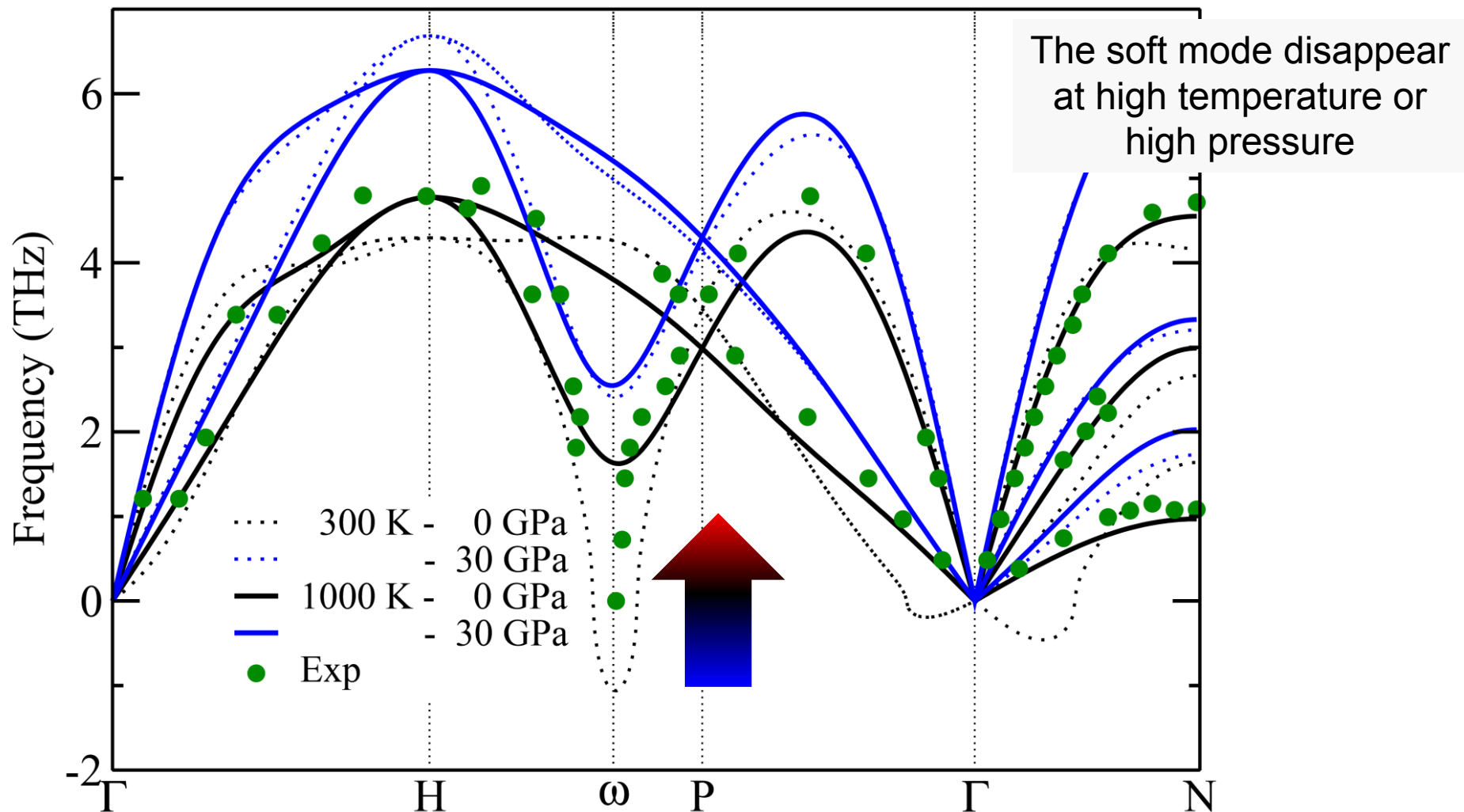


O. Hellman *et al.*, PRB **84**, 180301(R) (2011) ; O. Hellman *et al.*, PRB **87**, 104111 (2013).

J. Bouchet & F. Bottin., Phys. Rev. B **92**, 174108 (2015) ; F. Bottin, J. Bieder & J. Bouchet CPC **254**, 107301 (2020).



ZIRCONIUM (BCC) : STABILISATION AT HIGH T OR HIGH P

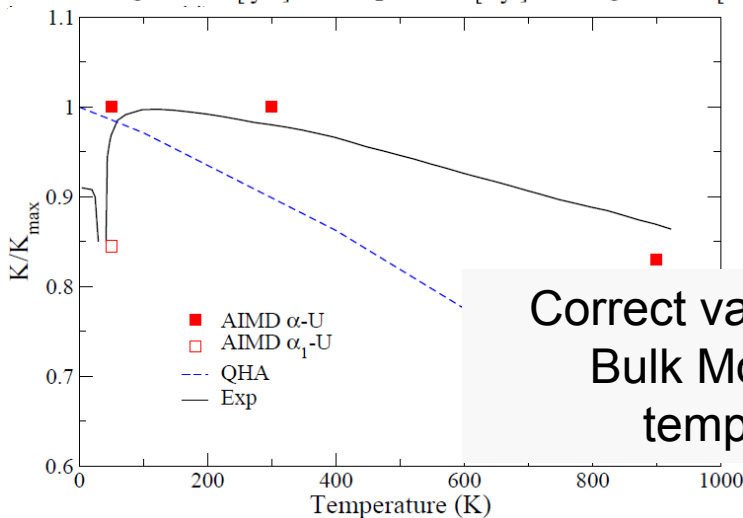
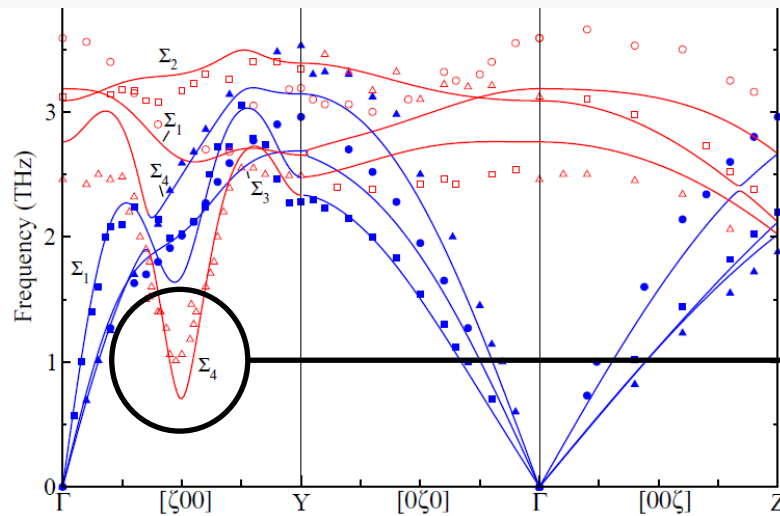


O. Hellman *et al.*, PRB **84**, 180301(R) (2011).

S. Anzellini, F. Bottin, J. Bouchet and A. Dewaele *et al.*, PRB **102**, 184105 (2020).

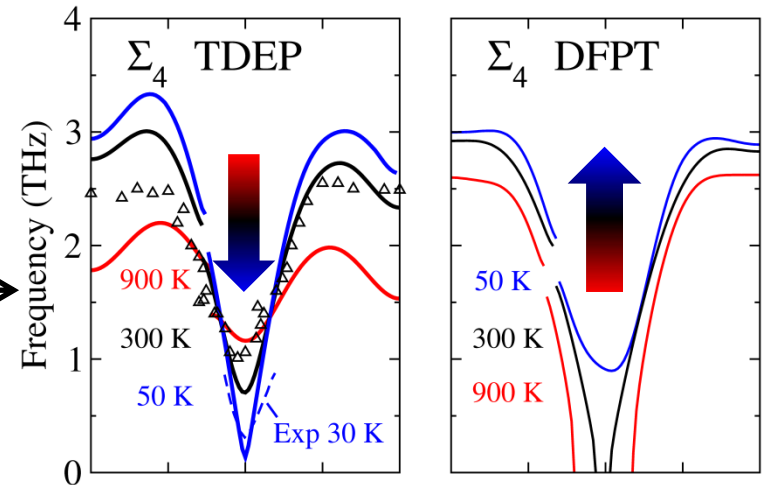
URANIUM (ALPHA) : FAILURE OF THE QHA

aTDEP phonon spectrum in good agreement with experiment (300K)

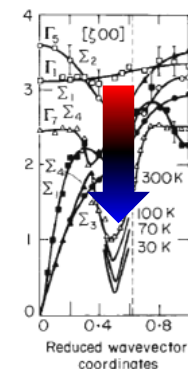


Correct variation of the Bulk Modulus wrt temperature

Correct variation of the phonon modes wrt temperature

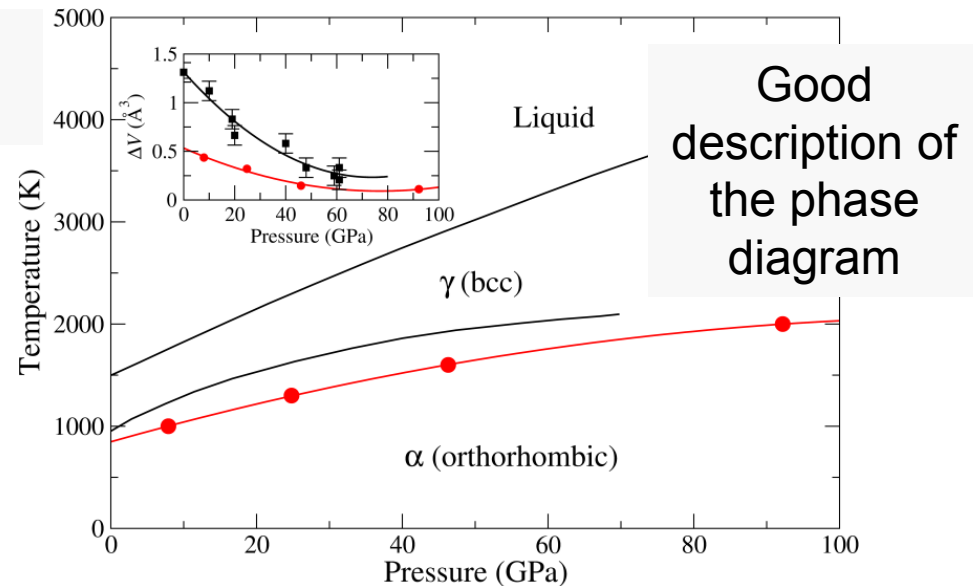
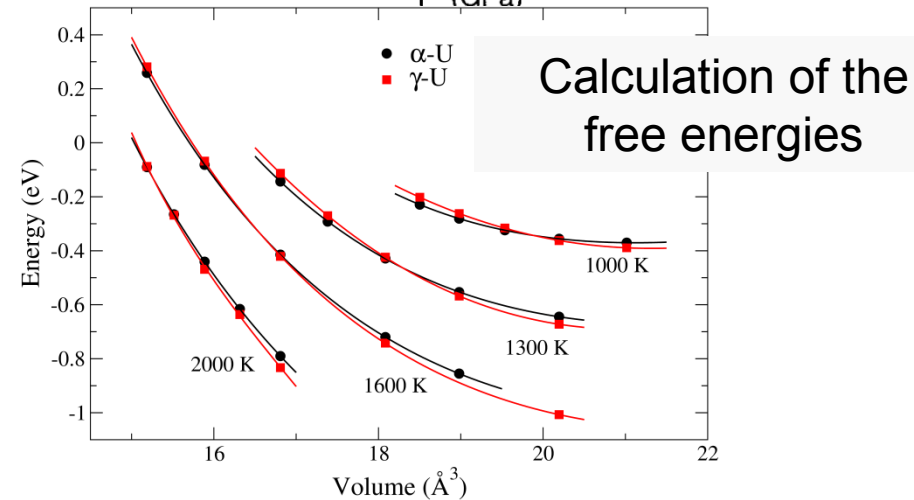
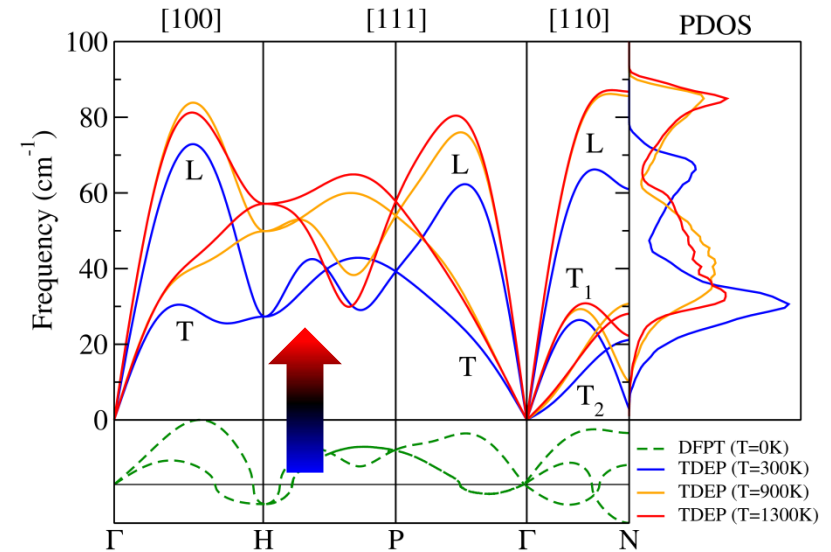
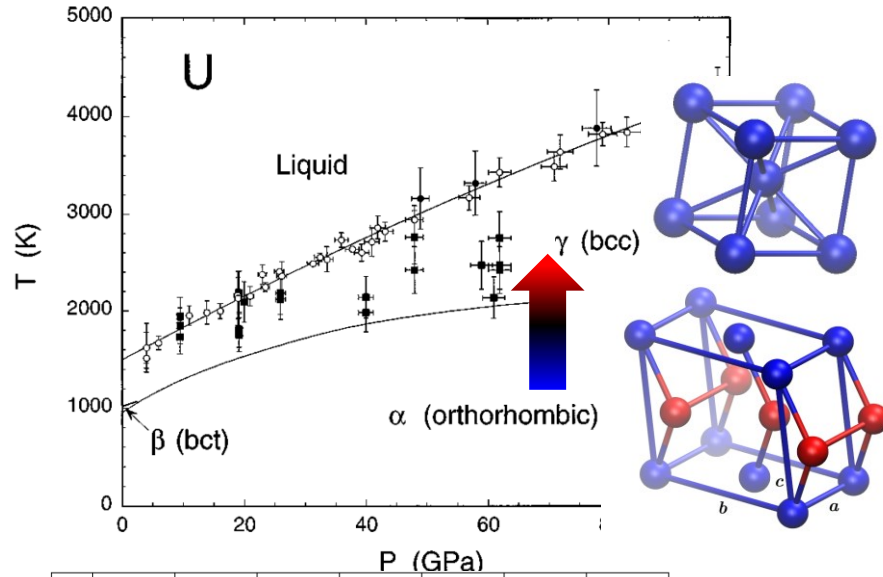


Exp



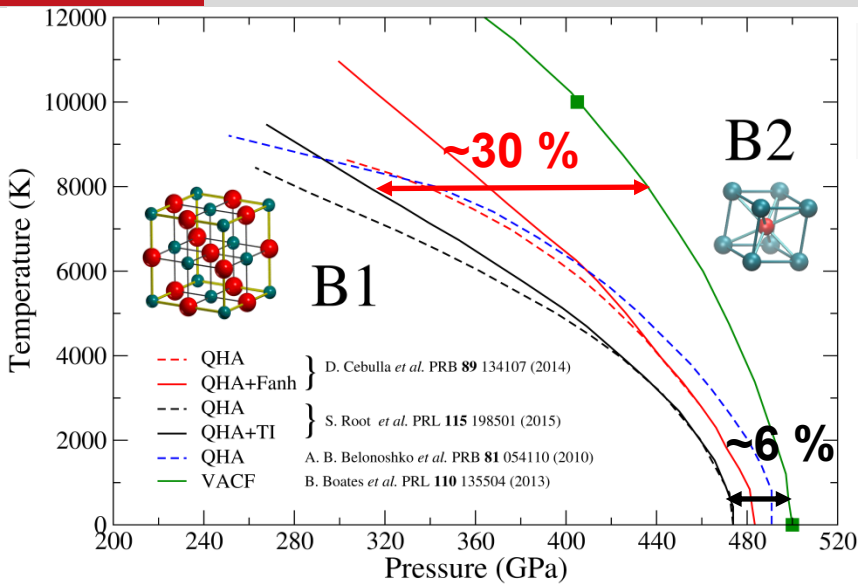
J. Bouchet & F. Bottin., Phys. Rev. B **92**, 174108 (2015)

URANIUM (GAMMA) : THE PHASE DIAGRAM

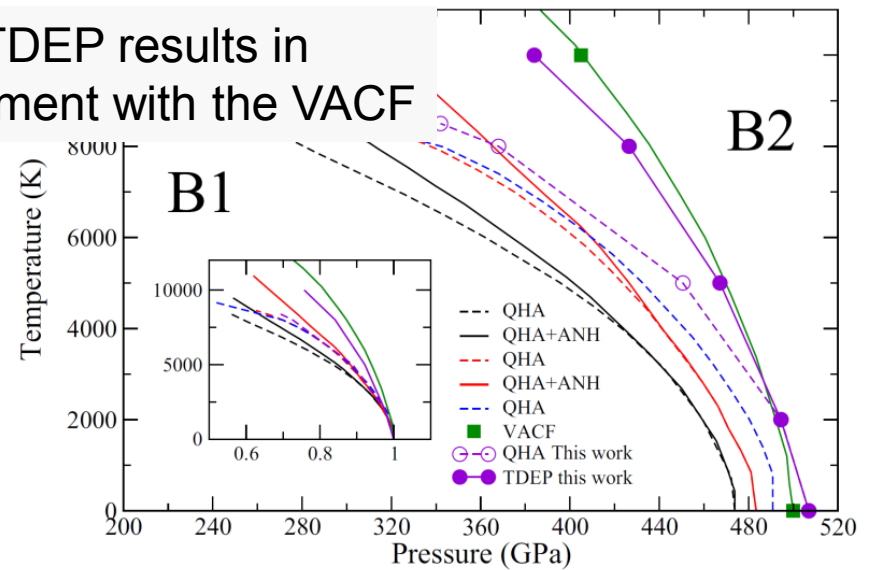


C.-S. Yoo, H. Cynn, & P. Söderlind PRB **57**, 10359 (1998)
J. Bouchet & F. Bottin., Phys. Rev. B **95**, 054113 (2017)

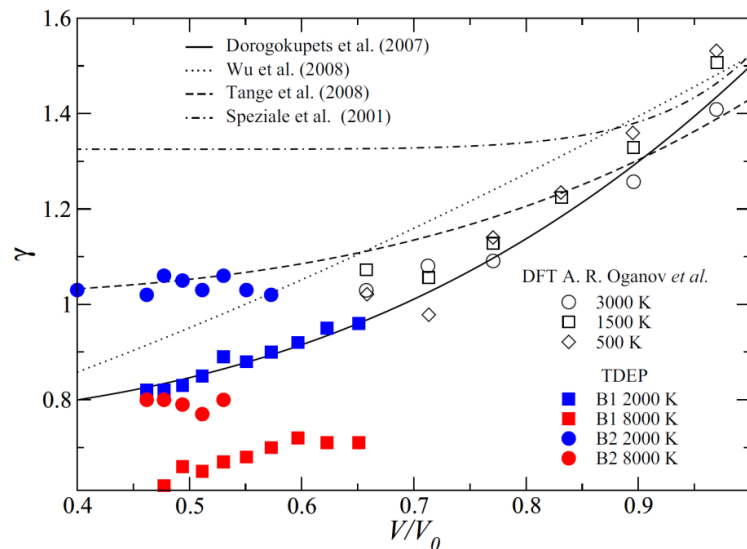
MGO : THE B1-B2 PHASE TRANSITION



TDEP results in agreement with the VACF



Variation of the Gruneisen parameter wrt temperature (along an isochoric process)



J. Bouchet, F. Bottin, V. Recoules, F. Remus, G. Morard, R.M. Bolis & A. Benuzzi-Mounaix, PRB 99, 094113 (2019)

As usual...

```
tdep < input.files > log
```

...with 3 lines in the input.files...

```
input.in  
HIST.nc  
output
```

... and a few variables in the input.in file (see the atdep suite) :

```
NormalMode  
# Unit cell definition  
    brav    7    0  
natom_unitcell    5  
xred_unitcell    0.0 0.0 0.0  0.5 0.5 0.5  0.0 0.5 0.5  0.5 0.0 0.5  0.5 0.5 0.0  
typat_unitcell    3    2    1    1    1  
# Supercell definition  
    multiplicity    2.00    0.00    0.00    0.00    2.00    0.00    0.00    0.00    2.00  
    temperature    495.05  
# Computation details  
    nstep_max        101  
    nstep_min         1  
    Rcut             7.426  
# Optional inputs  
    Ngqpt2 2 2 2  
TheEnd
```


A large number of files :

- Phonon frequencies (omega.dat)
- Dynamical matrices (dij.dat & DDB)
- Interatomic Force Constants (ifc.dat & ifc.nc)
- Thermodynamic data (thermo.dat)
- ...
- Main output file :

vibra#####

See the vdos.dat and TDEP Grun:#####
Write the IFC of TDEP in :#####

Grunelsen parameter, Thermal expansion, Thermal pressure... #####
See the grunelsen.dat, alpha_gamma.dat and thermo3.dat files

For each shell, list of coefficients (IFC), number of neighbours...

```
##### NEW SHELL (ishell= 1): There are 1 atoms on this shell
For latcell= 1 ,with type= 1
For jatom = 1 ,with type= 1
For katom = 1 ,with type= 1
For latom = 1 ,with type= 1
\Phi4^{ 1 1kl}=
0.315691 0.000000 0.000000
0.000000 0.105230 0.000000
0.000000 0.000000 0.114645
\Phi4^{ 1 2kl}=
0.000000 0.105230 0.000000
0.105230 0.000000 0.000000
0.000000 0.000000 0.000000
\Phi4^{ 1 3kl}=
0.000000 0.000000 0.114645
0.000000 0.000000 0.000000
0.114645 0.000000 0.000000
\Phi4^{ 2 1kl}=
0.000000 0.105230 0.000000
0.105230 0.000000 0.000000
0.000000 0.000000 0.000000
\Phi4^{ 2 2kl}=
0.105230 0.000000 0.000000
0.000000 0.315691 0.000000
0.000000 0.000000 0.114645
\Phi4^{ 2 3kl}=
0.000000 0.000000 0.000000
0.000000 0.000000 0.114645
0.000000 0.114645 0.000000
\Phi4^{ 3 1kl}=
0.000000 0.000000 0.114645
```

```
##### For each shell, list of coefficients (IFC), nu
##### NEW SHELL (ishell= 1): There are 1 ato
For latcell= 1 ,with type= 1
For jatom = 1 ,with type= 1
For katom = 1 ,with type= 1
\Phi3^{(\alpha\beta\gamma\delta)}=
0.000000 0.106589 0.000000
0.106589 0.000000 0.000000
0.000000 0.000000 0.000000
\Phi3^{(\alpha\beta\gamma\delta)}=
0.106589 0.000000 0.000000
0.000000 -0.106589 0.000000
0.000000 0.000000 0.000000
\Phi3^{(\alpha\beta\gamma\delta)}=
0.000000 0.000000 0.000000
0.000000 0.000000 0.000000
0.000000 0.000000 0.000000
(i,j) vector components: 0.000000 0.000000 0.000
(j,k) vector components: 0.000000 0.000000 0.000
(k,i) vector components: 0.000000 0.000000 0.000
```

```
##### NEW SHELL (ishell= 2): There are 18 ato
For latcell= 1 ,with type= 1
For jatom = 1 ,with type= 1
For katom = 2 ,with type= 2
\Phi3^{(\alpha\beta\gamma\delta)}=
0.007572 -0.015352 -0.047539
0.015352 0.002431 -0.042105
-0.047539 -0.042105 -0.096677
\Phi3^{(\alpha\beta\gamma\delta)}=
-0.015433 0.002571 -0.042105
0.002571 0.021208 0.001079
-0.042105 0.001079 -0.055817
(u,v,w) = (u_v-u_k)/(u_v+u_k) = 0.000
```

```
##### Elastic constant
##### Bulk and Shear modulus...
#####
===== Using the formulation proposed
Cijkl [in GPa]=
| C11 C12 C13 C14 C15 C16 | 105.595
| C21 C22 C23 C24 C25 C26 | 52.501
| C31 C32 C33 C34 C35 C36 | 52.501
| C41 C42 C43 C44 C45 C46 | 0.000
| C51 C52 C53 C54 C55 C56 | 0.000
| C61 C62 C63 C64 C65 C66 | 0.000

===== For an Anisotropic Material
Sijkl [in GPa-1]=
| S11 S12 S13 S14 S15 S16 | 0.014
| S21 S22 S23 S24 S25 S26 | -0.005
| S31 S32 S33 S34 S35 S36 | -0.005
| S41 S42 S43 S44 S45 S46 | 0.000
| S51 S52 S53 S54 S55 S56 | 0.000
| S61 S62 S63 S64 S65 S66 | 0.000

===== For an Orthotropic Material (S11, S12, S13, S14, S15, S16, S17, S18, S19, S20, S21, S22, S23, S24, S25, S26, S27, S28, S29, S30, S31, S32, S33, S34, S35, S36, S37, S38, S39, S40, S41, S42, S43, S44, S45, S46, S47, S48, S49, S50, S51, S52, S53, S54, S55, S56, S57, S58, S59, S60, S61, S62, S63, S64, S65, S66, S67, S68, S69, S70)
Young modulus E1, E2 and E3 [in GPa]= 70
Poisson ratio Nu21, Nu31, Nu23, Nu12, Nu13, Nu32
Shear modulus G23, G13 and G12 [in GPa]=
Sijkl [in GPa-1]=
| S11 S12 S13 S14 S15 S16 | 0.014
| S21 S22 S23 S24 S25 S26 | -0.005
| S31 S32 S33 S34 S35 S36 | -0.005
| S41 S42 S43 S44 S45 S46 | 0.000
| S51 S52 S53 S54 S55 S56 | 0.000
| S61 S62 S63 S64 S65 S66 | 0.000

===== Voigt
ISOTHERMAL modulus [in GPa]: B
Average of Young modulus E [in GPa]= 70
Velocities [in m.s-1]: compres
Debye velocity [in m.s-1]= 354

===== Reus
ISOTHERMAL modulus [in GPa]: B
Average of Young modulus E [in GPa]= 70
Velocities [in m.s-1]: compres
Debye velocity [in m.s-1]= 354

===== Elas
Elastic anisotropy index : A_U
Bulk anisotropy ratio : A_B= (B/B0)
Shear anisotropy ratio : A_G= (G/G0) = 0.000
```

- ✓ **Implementation of the :**
 - **4th order IFCs**
 - **Dipole-dipole interaction (LO-TO splitting)**
 - **Alloy description (VCA)**
 - **Reweighting (see the presentation of A. Castellano)**
 - **Parallelisation over configurations**

- ✓ **Creation of an aTDEP series of test (around 40 tests)**
- ✓ **Improvement of the documentation for the user (see the « Topic », the « Input variables » and the « User guide »)**
- ✓ **Building of a first « aTDEP tutorial »**

ATDEP : USER GUIDE – TOPIC – INPUT VARIABLES - TUTORIAL

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Spacegroup

Installation

The aTDEP utility

The Temperature Dependent Effective Potential (TDEP) method has been developed by O. Hellman et al. [Hellman2011], [Hellman2013], [Hellman2013a] in 2011 and the aTDEP implementation in ABINIT has been performed and used for the first time in 2015 by J. Bouchet and F. Bottin [Bouchet2015], [Bouchet2017].

- User guide (in a pdf format): [aTDEP guide](#)
- Theory: [aTDEP paper](#) corresponding to the article [Bottin2020]

Prerequisite and theory

The approach used in this code is detailed in a publication dedicated to the development of all formula (see [aTDEP paper](#)). We strongly encourage all the users to carefully read this paper before beginning. All the vibrational, elastic and thermodynamic quantities computed by aTDEP are presented with the same writing conventions as the ones used in the output files of aTDEP. In the same manner, a comprehensive understanding of some ABINIT basic variables is also required in order to fill the input file and read the output file of aTDEP.

In addition, this paper is also useful to understand the limitations and convergences which are inherent to the present method. These particular points are sometimes discussed in the article, with some references and illustrating examples.

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- Prerequisite and theory
- The ABINIT computation
- The aTDEP computation
- The input files
- The output files

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Topics

Features

Topic List

Abipy

APPA

Artificial

AtomManipulator

aTDEP

AtomTypes

Bader

Band2Eps

Berry

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BSE

BoundingProcess

ConstrainedDFT

ConstrainedPol

Control

Coulomb

CrossingBarriers

CRPA

Crystal

aTDEP

This page gives hints on how to perform thermodynamic, elastic and transport properties calculations including explicit temperature effects with the ABINIT package.

- User guide: [aTDEP guide](#)
- Theory: [aTDEP paper](#) corresponding to the article [Bottin2020]

Introduction

The Temperature Dependent Effective Potential (TDEP) method has been developed by O. Hellman et al. [Hellman2011], [Hellman2013], [Hellman2013a] in 2011 and the [aTDEP] implementation in ABINIT has been performed and used for the first time in 2015 by J. Bouchet and F. Bottin [Bouchet2015], [Bouchet2017].

The capture of thermal effects in solid state physics is a long standing issue and several stand-alone or post-process computational codes are available. Using different theoretical frameworks, they propose to provide some thermodynamic quantities involving the so called anharmonic effects. [aTDEP] calculation can produce almost all the temperature-dependent thermodynamic quantities you want, from a single *ab initio* molecular dynamic (AIMD) trajectory and by means of a Graphical User Interface (GUI) very easy to use (AGATE).

The original TDEP method [Hellman2011] is implemented in ABINIT. In particular, various

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- Introduction
- Related Input Variables
- Selected Input Files

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Variables

All Variables

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External Parameters

Statistics

atdep input variables

This document lists and provides the description of the name (keywords) of the atdep input variables to be used in the input file for the atdep executable.

alloy

Mnemonics: ALLOY treatment
Mentioned in topic(s): [topic_aTDEP](#)
Variable type: integer
Dimensions: 3
Default value: 0
Added in version: 9.5.1

[Test list \(click to open\). Moderately used, \[1/43\] in all atdep tests, \[0/5\] in atdep tutorials](#)

OPTIONAL: Defines the treatment of the alloy. The first value defines the approximation used (0=nothing and 1=Virtual Crystal Approximation). The second and the third values define the [[typat_unitcell]] of the atoms which have to be alloyed.

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- alloy
- amu
- angle
- born_charge
- brav
- bzlength
- bzpath
- dielec_constant
- dosdeltae
- enunit
- multiplicity
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- natom_unitcell
- ngqpt1
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- nstep_max
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The 2nd order effective Interatomic Force Constants