Electron-phonon beyond Fröhlich : dynamical quadrupoles in polar and covalent solids ABIDEV 2021

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Electron-phonon coupling governs many phenomena



[Yu & Cardona, Fundamentals of Semiconductors]



[Margine, EPW school 2018]

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In semiconductors, at high T or when the impurity concentration is very low, the mobility is governed by electron-phonon interactions.

Electron mobility in the self-energy relaxation-time approximation to the Boltzmann transport equation :

$$\mu_{e,\alpha\beta} = \frac{-1}{\Omega n_e} \sum_{n \in CB} \int \left. \frac{d\mathbf{k}}{\Omega_{BZ}} \mathbf{v}_{n\mathbf{k},\alpha} \mathbf{v}_{n\mathbf{k},\beta} \tau_{n\mathbf{k}} \left. \frac{\partial f^0}{\partial \varepsilon} \right|_{\varepsilon_{n\mathbf{k}}} \right.$$

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Electronic lifetimes due to the scattering by phonons :

$$\frac{1}{\tau_{n\mathbf{k}}} = \frac{2\pi}{\hbar} \sum_{m,\nu} \int_{\mathsf{BZ}} \frac{d\mathbf{q}}{\Omega_{\mathsf{BZ}}} |g_{mn\nu}(\mathbf{k},\mathbf{q})|^2 \times [(n_{\nu\mathbf{q}} + f_{m\mathbf{k}+\mathbf{q}})\delta(\varepsilon_{n\mathbf{k}} - \varepsilon_{m\mathbf{k}+\mathbf{q}} + \omega_{\mathbf{q}\nu}) + (n_{\nu\mathbf{q}} + 1 - f_{m\mathbf{k}+\mathbf{q}})\delta(\varepsilon_{n\mathbf{k}} - \varepsilon_{m\mathbf{k}+\mathbf{q}} - \omega_{\mathbf{q}\nu})]$$

$$\begin{split} \frac{1}{\tau_{n\mathbf{k}}} &= 2\pi \sum_{m,\nu} \int_{\mathsf{BZ}} \frac{d\mathbf{q}}{\Omega_{\mathsf{BZ}}} |g_{mn\nu}(\mathbf{k},\mathbf{q})|^2 \times \left[(n_{\mathbf{q}\nu}^0 + f_{m\mathbf{k}+\mathbf{q}}^0) \delta(\varepsilon_{n\mathbf{k}} - \varepsilon_{m\mathbf{k}+\mathbf{q}} + \omega_{\mathbf{q}\nu}) \right. \\ &+ (n_{\mathbf{q}\nu}^0 + 1 - f_{m\mathbf{k}+\mathbf{q}}^0) \delta(\varepsilon_{n\mathbf{k}} - \varepsilon_{m\mathbf{k}+\mathbf{q}} - \omega_{\mathbf{q}\nu}) \right] \\ g_{mn\nu}(\mathbf{k},\mathbf{q}) &= \text{electron-phonon coupling matrix elements} = \langle \psi_{m\mathbf{k}+\mathbf{q}}^{\mathsf{KS}} | \Delta_{\mathbf{q}\nu} V^{\mathsf{KS}} | \psi_{n\mathbf{k}}^{\mathsf{KS}} \rangle. \end{split}$$

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Problem : this integral converges very slowly ($\sim 100 \times 100 \times 100$ q mesh required).

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Problem : this integral converges very slowly (~ $100 \times 100 \times 100$ q mesh required). **Solution :** Wannier interpolation of the $g_{mn\nu}(\mathbf{k}, \mathbf{q})$ obtained from DFPT $\Rightarrow EPW$ approach.



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Problem : this integral converges very slowly ($\sim 100 \times 100 \times 100 \text{ q}$ mesh required). **Our solution :** Fourier interpolation of the $\Delta_{q\nu} V^{KS}$ obtained from DFPT, and computation of ψ_{mk+q}^{KS} .



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[PRL 115, 176401 (2015), PRB 92, 054307 (2015)]

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We need to consider different important interactions for a correct physical interpretation :

- Dipoles are not the only electrostatic effect appearing when atoms are displaced.
- Dynamical quadrupoles are also created, even in nonpolar semiconductors.



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The scattering potentials can be written in the phonon mode ($\mathbf{q}\nu$) or atomic representation ($\mathbf{q}\kappa\alpha$) :

$$\Delta_{\mathbf{q}
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The interpolation procedure is similar to what is done for the dynamical matrix to obtain phonon spectra :

- Split the DFPT scattering potentials into long-range (V^L) and short-range (V^S) parts.
- **2** Interpolate V^S only.
- **3** Add $V^{\mathcal{L}}$ to the interpolated $V^{\mathcal{S}}$.

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$$V^{\mathcal{L}}_{_{\kappa\alpha,\mathbf{q}}}(\mathbf{r})=V^{D}_{_{\kappa\alpha,\mathbf{q}}}(\mathbf{r})+V^{Q}_{_{\kappa\alpha,\mathbf{q}}}(\mathbf{r})+V^{\mathcal{E}}_{_{\kappa\alpha,\mathbf{q}}}(\mathbf{r})$$

- **1** Split the DFPT scattering potentials into long-range $(V^{\mathcal{L}})$ and short-range $(V^{\mathcal{S}})$ parts.
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$$V^{\mathcal{L}}_{\kappalpha,\mathbf{q}}(\mathbf{r}) = V^{D}_{\kappalpha,\mathbf{q}}(\mathbf{r}) + V^{Q}_{\kappalpha,\mathbf{q}}(\mathbf{r}) + V^{\mathcal{E}}_{\kappalpha,\mathbf{q}}(\mathbf{r})$$

$$egin{aligned} V^D_{\kappalpha, \mathbf{q}}(\mathbf{r}) \propto rac{\mathbf{q}\cdot\mathbf{Z}^*}{\mathbf{q}\cdot\epsilon^\infty\cdot\mathbf{q}} \ V^Q_{\kappalpha, \mathbf{q}}(\mathbf{r}) \propto rac{\mathbf{q}\cdot\mathbf{Q}^*\cdot\mathbf{q}}{\mathbf{q}\cdot\epsilon^\infty\cdot\mathbf{q}} \end{aligned}$$

Interpolating the scattering potentials without the dynamical quadrupoles interactions is bad



- We start from a given N_q × N_q × N_q DFPT
 q-point grid and interpolate the scattering potentials on the fine meshes required to converge the mobility.
- Without including the dynamical quadrupoles, the interpolation converges slowly.
- Including them fixes the interpolation and the correct mobility can be obtained.

In-depth analysis : which modes are the most impacted by quadrupoles ?

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- Quadrupole acoustic sum rule for nonpolar semiconductors ⇒ larger impact on optical modes.
- Piezoelectric coupling in polar materials ⇒ larger impact on acoustic modes.

Dynamical quadrupoles seem particularly important for strongly piezoelectric systems

Example : ZnO



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- compute the electron and/or hole phonon-limited mobility starting from a DFPT job (DDB, DVDB files) and a NSCF run
 - possibility to greatly reduce the NSCF part by computing only the k points close to the CBM/VBM using the "kerange" trick
 - **k** and **q** meshes up to 400×400×400 can be reached!

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- use robots to compare the results obtained by different runs (convergence of the mobility, comparison of linewidths,...),
- use a workflow to compute the mobility directly : all that is needed is the structure, pseudopotentials, and the size of the meshes, see abipy/examples/flows/run_eph_mob.py

- The implementation in ABINIT to compute the phonon-limited lifetimes (eph_task -4) and transport tensors (eph_task -4 or 7) is ready and documented with a tutorial.
- Computations are rather optimized (many parallelization levels) and can be run through an AbiPy workflow.
- Dynamical quadrupoles are required for a correct interpolation of the e-ph scattering potentials. They can be computed with ABINIT as well! Unphysical results can be obtained otherwise.
- Check our papers : Phys. Rev. Lett. 125, 136601 (2020) and Phys. Rev. B 102, 094308 (2020).

Thank you!