DE LA RECHERCHE À L'INDUSTRIE



#### HYBRID FUNCTIONALS IN ABINIT:

# STATE OF THE ART AND PERSPECTIVES

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#### Implementation of Fock-exchange with PAW

**Energy** 

**Forces** 

**Stresses** 

Range-separated hybrid functionals

Perspectives: Performance of the calculation- the ACE method

Implementation of Fock-exchange with PAW

# Fock exchange term:

$$E_{exch}^{\uparrow} = -\frac{1}{2} \sum_{\mu}^{val} \sum_{\nu}^{val} f_{\mu} f_{\nu} \delta_{\sigma_{\mu}, \sigma_{\nu}} \left( \left( \tilde{n}_{\mu\nu} + \hat{n}_{\mu\nu} \right) \right) - \frac{1}{2} \sum_{\sigma} \sum_{ijkl} \rho_{ij}^{*} (\sigma) \rho_{kl} (\sigma) e_{ikjl} - \sum_{ij} \rho_{ij} X_{ij} + E_{x}^{c-c}$$

with 
$$\begin{split} & \left( (n_{\mu\nu}) \right) = \int d\mathbf{r} d\mathbf{r}' \frac{n_{\mu\nu}^*(r) n_{\mu\nu}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} \\ & n_{\mu\nu}(\mathbf{r}) = \psi_{\mu}^*(r) \psi_{\nu}(r) \\ & \hat{n}_{\mu\nu}(\mathbf{r}) = \sum_{LM} \sum_{a} \sum_{i,j} Q_{i,j}^{a,LM}(\mathbf{r}) \; \langle \tilde{\psi}_{\mu} | \tilde{p}_{i}^{a} \rangle \; \langle \tilde{p}_{j}^{a} | \tilde{\psi}_{\nu} \rangle \\ & X_{ij}^{a} = \frac{1}{2} \sum_{c} \iint \frac{\phi_{i}(r) \phi_{j}(r') \phi_{c}(r) \phi_{c}(r')}{|r - r'|} dr dr' \\ & E_{x}^{c-c} \end{split}$$
 given in the JTHv1.0 PAW atomic data files

eikjl is already calculated for the GS

#### Fock exchange: the new Hamiltonian

By considering Fock exchange, the Hamiltonian becomes:

$$\widetilde{H}|\widetilde{\psi}_{\nu}\rangle = \left[-\frac{1}{2}\nabla^{2} + \widetilde{v}_{\text{eff}}(\mathbf{r})\right]|\widetilde{\psi}_{\nu}\rangle + \sum_{a}\sum_{i,j}|\widetilde{p}_{i}^{a}\rangle D_{ij}^{a}\langle\widetilde{p}_{j}^{a}|\widetilde{\psi}_{\nu}\rangle$$

$$+\sum_{\mu \text{ occ.}} f_{\mu} \left[ v_x^F(\mathbf{r}) \right]_{\mu\nu} \left| \widetilde{\psi}_{\mu} \right\rangle$$

$$+\sum_{\mu \text{ occ.}} f_{\mu} \left[v_{x}^{F}(\mathbf{r})\right]_{\mu\nu} \left|\widetilde{\psi}_{\mu}\right\rangle + \sum_{a} \sum_{i,j} \left|\widetilde{p}_{i}^{a}\right\rangle D_{ij}^{Fock \ a} \left\langle \widetilde{p}_{j}^{a} \middle|\widetilde{\psi}_{\nu}\right\rangle + \sum_{a} \sum_{i,j} \sum_{\mu \text{ occ.}} f_{\mu} \left|\widetilde{p}_{i}^{a}\right\rangle \hat{D}_{ij}^{Fock \ a}(\mu,\nu) \left\langle \widetilde{p}_{j}^{a} \middle|\widetilde{\psi}_{\mu}\right\rangle$$

The first term is the well-known Fock-exchange potential:

$$[v_x^F(\mathbf{r})]_{\mu\nu} = \int d\mathbf{r}' \frac{\widetilde{n}_{\mu\nu}(\mathbf{r}') + \widehat{n}_{\mu\nu}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

The second term is a standard non-local Dij term:

$$D_{ij}^{Fock\,a} = -\sum_{kl} \rho_{kl} e_{ikjl} - X_{ij}^{a}$$

The last term is a non-standard non-local Dij term:

A structure similar as the usual

$$\hat{D}_{ij}^{Fock\ a}(\mu,\nu) = -\sum_{LM} \int d\mathbf{r} \ [v_x^F(\mathbf{r})]_{\mu\nu} \ Q_{ij}^{a,LM}(\mathbf{r})$$

This term can be calculated with the routine paw dijhat BUT on the fly in fock getghc.



#### **Fock forces**

- Norm-conserving case: no contribution to the forces
- PAW case: two new terms

$$\sum_{\nu} \sum_{\mu occ} \sum_{a} \sum_{ij} f_{\nu} f_{\mu} \left\langle \tilde{\psi}_{\nu} \frac{\partial}{\partial \mathbf{R}} ( | \tilde{p}_{i}^{a} \rangle \left\langle \tilde{p}_{j}^{a}) \middle| \tilde{\psi}_{\mu} \right\rangle \widehat{D}_{ij}^{fock \ a} (\mu, \nu) \qquad \text{Calculated in nonlop.F90}$$

$$\sum_{\nu} \sum_{\mu occ} f_{\mu} f_{\nu} \int [v_{x}^{F}(\mathbf{r})]_{\mu\nu} \frac{\partial \hat{n}_{\mu\nu}(\mathbf{r})}{\partial (\mathbf{r} - \mathbf{R})} d\mathbf{r}$$
 Calculated in pawmknhat\_psipsi.F90



Coded and tested in ABINIT

#### **Fock stresses**

■ Norm-conserving case:

$$strfock(\tilde{n}_{\mu\nu}) = -\frac{\delta_{\alpha\beta} E_{\mu\nu}^{Fock}}{\Omega} - \sum_{G \neq 0} 4\pi \frac{G_{\alpha} G_{\beta}}{G^4} \tilde{n}_{\mu\nu} (G) \tilde{n}_{\mu\nu} (-G)$$



Equivalent to harstr in ABINIT and calculated in strfock.F90

Coded and tested in ABINIT

☐ PAW case:

$$\sigma_{\alpha\beta} = \sum_{\mu\nu} f_{\mu} f_{\nu} strfock (\tilde{n}_{\mu\nu} + \hat{n}_{\mu\nu}) + \int [v_{xc}^{F}(r)]_{\mu\nu} \left( \delta_{\alpha\beta} \, \hat{n}_{\mu\nu} + \frac{\partial \hat{n}_{\mu\nu}}{\partial \varepsilon_{\alpha\beta}} \right) d\mathbf{r}$$

$$\sum_{\nu} \sum_{\mu o c c} \sum_{a} \sum_{ij} f_{\nu} f_{\mu} \left\langle \tilde{\psi}_{\nu} \frac{\partial}{\partial \varepsilon_{\alpha\beta}} (|\tilde{p}_{i}^{a}\rangle \langle \tilde{p}_{j}^{a}) |\tilde{\psi}_{\mu}\rangle \widehat{D}_{ij}^{fock\ a}(\mu, \nu) \right\rangle$$
 Calculated in nonlop.F90

$$-\sum_{\nu}\sum_{\mu occ}f_{\mu}f_{\nu}\int[v_{x}^{F}(\mathbf{r})]_{\mu\nu}(\mathbf{r}-R)_{\beta}\frac{\partial\hat{n}_{\mu\nu}(\mathbf{r})}{\partial(\mathbf{r}-\mathbf{R})_{\alpha}}d\mathbf{r}$$

Coded in ABINIT (but bug...)

# Special attention to the norm-conserving case

$$V_x^{hyb}\left[n_c, n_v\right] = V_x^{DFT}\left[n_c + n_v\right] - \alpha V_x^{DFT}\left[n_v\right] + \alpha V_x^{Fock}\left[n_v\right]$$

But libxc gives:

$$V_x^{Libxc}[n] = (1 - \alpha)V_x^{DFT}[n]$$

Three calls to libxc are therefore necessary:

ee calls to libxc are therefore necessary: 
$$V_x^{hyb}\left[n_v\right]=(1-\alpha)V_x^{DFT}\left[n_v\right] \qquad \qquad (1)$$
 
$$V_x^{DFT}\left[n_c+n_v\right] \qquad \qquad (2)$$
 
$$V_x^{DFT}\left[n_v\right] \qquad \qquad (3)$$



done in xchybrid\_ncpp\_cc.F90 for energy, forces and stresses

# Range-separated hybrid functionals

$$V_{XC}^{HSE} = \alpha V_X^{Fock,SR}(\omega) + (1 - \alpha)V_X^{PBE,SR}(\omega) + V_X^{PBE,LR}(\omega) + V_C^{PBE}$$

$$\frac{1}{r} = \frac{erfc(\omega r)}{r} + \frac{erf(\omega r)}{r} = SR(\omega) + LR(\omega)$$

- lacktriangle Norm-conserving case: the erfc screening is taken into account in the calculation of  $[v_{xc}^F(r)]_{\mu\nu}$
- □ PAW case: some integrals contributing to the eijkl term must be screened with the erfc function.

This requires the development: 
$$\frac{erfc(\omega|\mathbf{r}-\mathbf{r}'|)}{|\mathbf{r}-\mathbf{r}'|} = \sum_{l=\grave{\mathbf{a}}}^{\infty} \omega \Phi_l(\omega R_>, \omega r_<) P_l(\widehat{\cos(\mathbf{r}, \mathbf{r}')})$$

This is done in poisson.F90 and screened\_coul\_kernel.F90



HSE03 and HSE06 can be achieved with ixc=-427 and -428



# How to calculate the Fock-exchange in practice

- For an ABINIT user, to make a calculation of Fock exchange:
  - Do a first dataset for Ground State
  - Do a second dataset for Fock calculation choosing ixc=40-42 (HF, PBE0, PBE0-1/3)

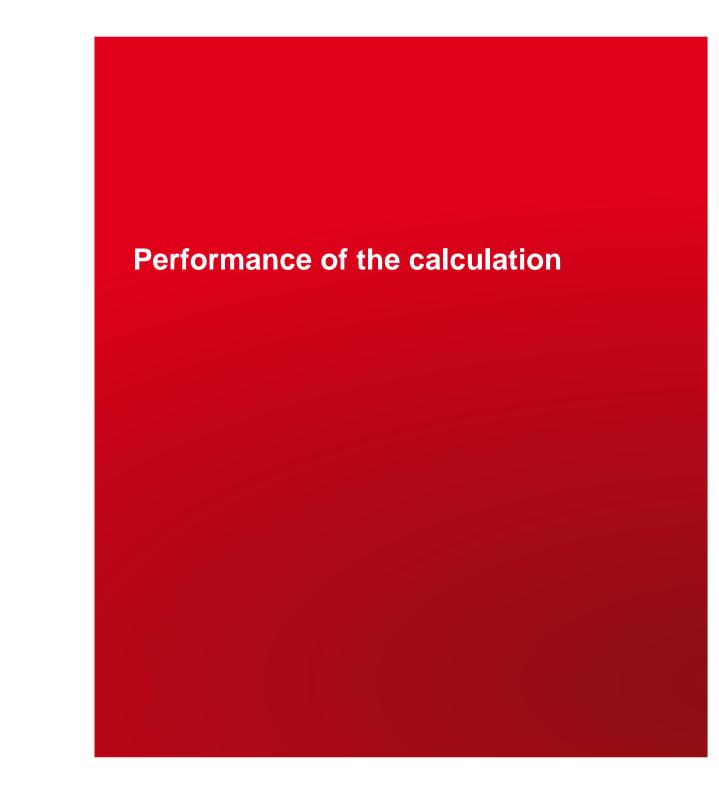
=-406 (PBE0-Libxc)

= -456 (PBE0-1/3-Libxc)

=-427,-428 (HSE03,HSE06)

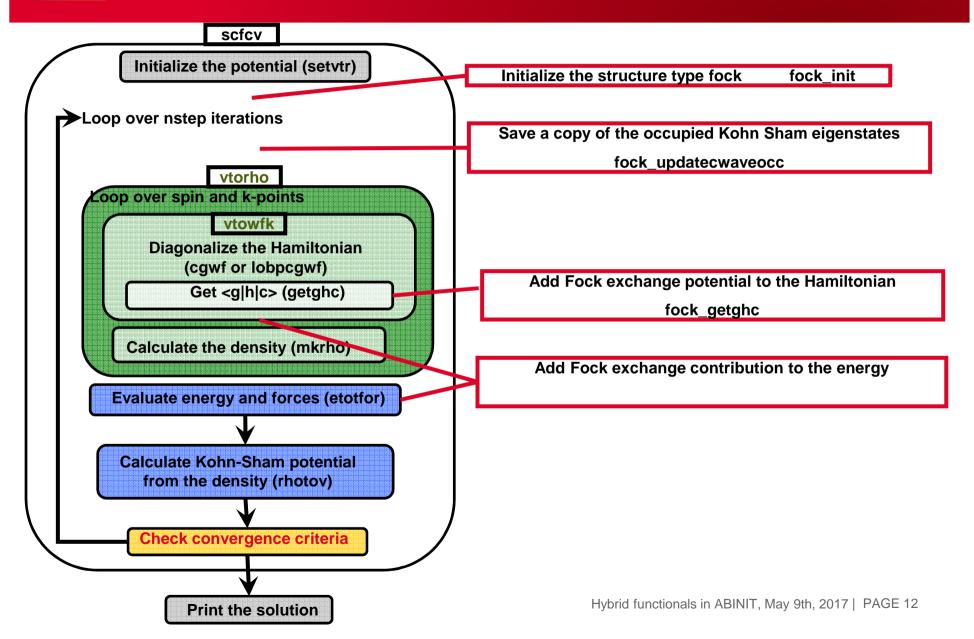
- npkpt : number of processors for k-point parallelization
- nphf: number of processors for Fock parallelization
- nbandhf: max. number of occupied valence states
- nnsclohf: number of loop without updating the occupied states cwaveocc default value: 1 (update at each iteration) in practice, 3 seems a good number.

WARNING: the spin polarized case has not been extensively tested use istwfk=1; iscf=2; paral\_kgb=0, paral\_atom=0 do not use simultaneously optstress and optforces



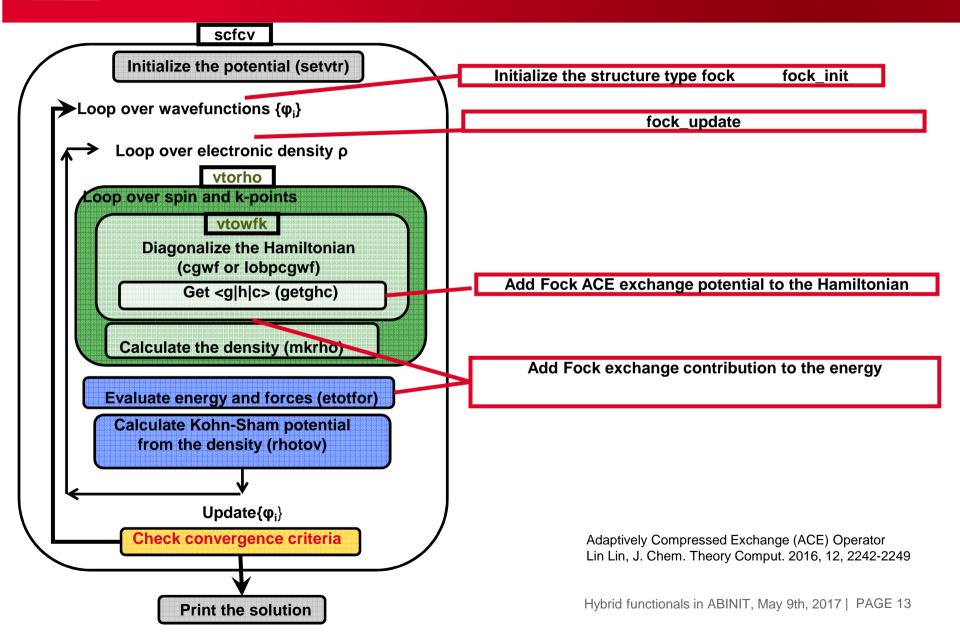


# **Future self consistent cycle**





# Future self consistent cycle: the ACE method



# New self consistent cycle: the ACE method

#### Fock update:

Calculate 
$$[v_{xc}^F(r)]_{\mu\nu}$$

Calculate 
$$W_{\nu}(r) = \sum_{\mu occ} f_{\mu} \left[ v_{xc}^F(r) \right]_{\mu \nu} \left| \tilde{\psi}_{\nu} > \right|$$

Calculate 
$$M_{kl} = \int \tilde{\psi}_k^*(\mathbf{r}) W_l(\mathbf{r}) d\mathbf{r} = -L_{kl} L_{kl}^T$$

Calculate 
$$\xi_k(r) = \sum_i W_i(r) (L^{-T})_{ik}$$

Calculate 
$$V_X^{ACE}(\mathbf{r}, \mathbf{r}') = -\sum_k \xi_k^*(\mathbf{r}) \xi_k(\mathbf{r}')$$
 Fock operator

Advantage: the full Fock term is calculated only for the loop on the wavefunctions.

The calculation of the ACE Fock term for the loop on the density costs the price of a non-local operator

