

10th ABINIT International Developer Workshop

May 31 - June 4, 2021

Are plane-waves basis packages able to predict the correct long range behavior of the local kinetic energy per particle?

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Camilo Espejo, Hugo Bohórquez and Aldo Romero





Summary

1. Background and Motivation.
2. Two theorems involving $n(\mathbf{r})$ and $G(\mathbf{r})$ in finite systems.
3. Computing the asymptotic behavior of $G(\mathbf{r})$ in **ABINIT** for atoms and molecules.
4. Conclusions and perspectives.

Background and Motivation

nature
materials

ARTICLES

<https://doi.org/10.1038/s41563-017-0012-2>

X-ray electron density investigation of chemical bonding in van der Waals materials

Hidetaka Kasai^{1,2}, Kasper Tolborg², Mattia Sist², Jiawei Zhang², Venkatesha R. Hathwar¹, Mette Ø. Filsø², Simone Cenedese², Kunihisa Sugimoto³, Jacob Overgaard², Eiji Nishibori¹ and Bo B. Iversen^{2*}

Nature Materials 17:249 (2018)

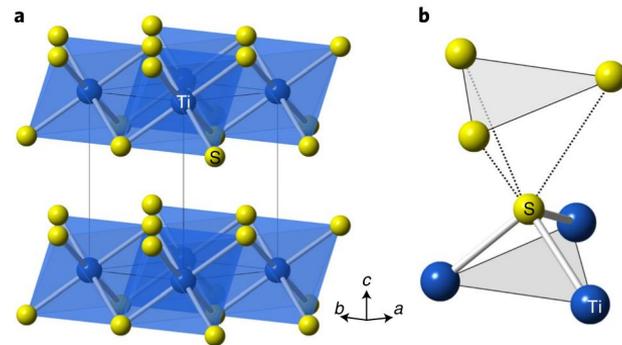


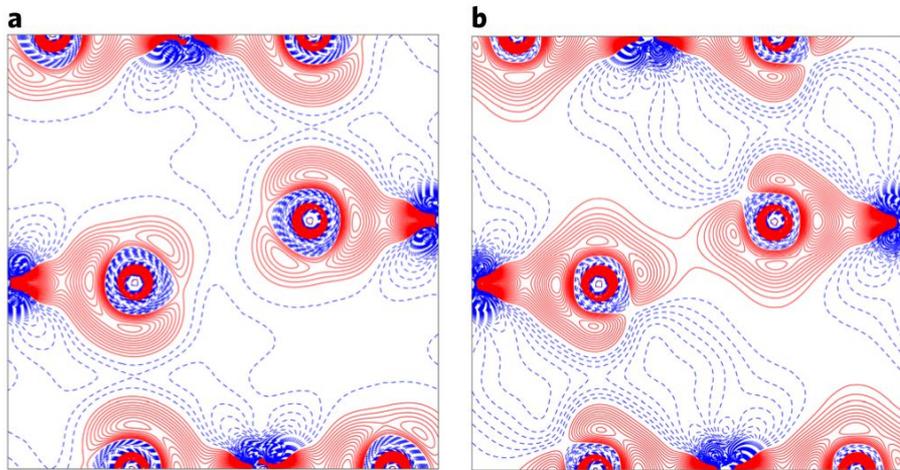
Fig. 1 | Structure of 1T-TiS₂. **a**, Ti atom is at the edge of a unit cell (Space group: $P\bar{3}m1$, $a = 3.3976(3)$ and $c = 5.6654(3)$ Å) forming a TiS₆ octahedron. **b**, S atom and its adjacent atoms in the direction of the c -axis.

Background and Motivation

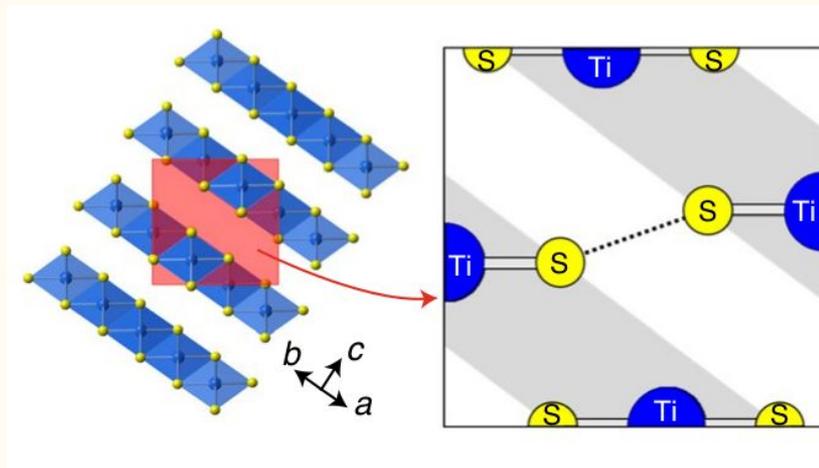


Experiment

Theory(SCAN+rVV10)



interaction. a,b, Static deformation map on the Ti-S-S-Ti plane (110 plane) from multipole modelling of the theoretical (SCAN+rVV10) (a) and experimental (b) structure factors. The contour interval is $0.01 e \text{ \AA}^{-3}$, with positive and negative contours drawn as solid red and dotted blue lines, respectively. **c,d**, Negative Laplacian map obtained



Even though both equilibrium interlayer distance and binding energy follow the experimental trend, there are important differences in the ground state electronic density.

Background and motivation



Table 2 | Atomic properties of Ti and S

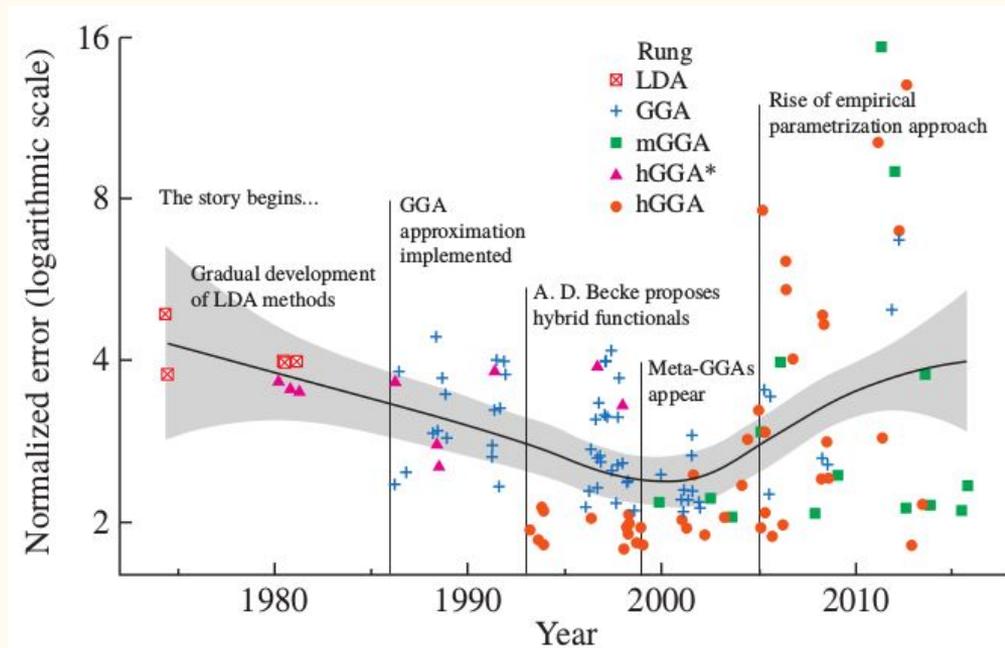
Variable	Method	Ti	S
Q	Experiment	1.63	-0.82
	Theory (SCAN+rVV10)	1.58	-0.80
V (Å ³)	Experiment	9.89	23.47
	Theory (SCAN+rVV10)	9.64	23.63
μ (eÅ)	Experiment	0.00	0.03
	Theory (SCAN+rVV10)	0.00	0.34
L (a.u.)	Experiment	-7.47×10^{-3}	-4.34×10^{-4}
	Theory (SCAN+rVV10)	-1.05×10^{-2}	-1.36×10^{-4}

Q, V and μ are the atomic charge, the atomic basin volume and the magnitude of atomic dipole moment, respectively. L is the atomic Lagrangian $(-(1/4) \int_{\Omega} \nabla^2 \rho \, d\tau)$, which is a measure of the accuracy of the integration.

Nature Materials 17:249 (2018)

Houston, we have a problem!

Medvedev, *et. al. Science* **355**, 49-52(2017)



REPORT

THEORETICAL CHEMISTRY

Density functional theory is straying from the path toward the exact functional

Michael G. Medvedev,^{1,2,3*†} Ivan S. Bushmarinov,^{1*†} Jianwei Sun,^{4‡} John P. Perdew,^{4,5†} Konstantin A. Lyssenko^{1†}

Mendelev Commun., 2018, 28, 225–235

“The reproduction of the exact energy is not a characteristic of the exact functional unless the input electronic density is exact as well.”

Houston, we have a problem!

Medvedev, *et. al. Science* **355**, 49-52(2017)

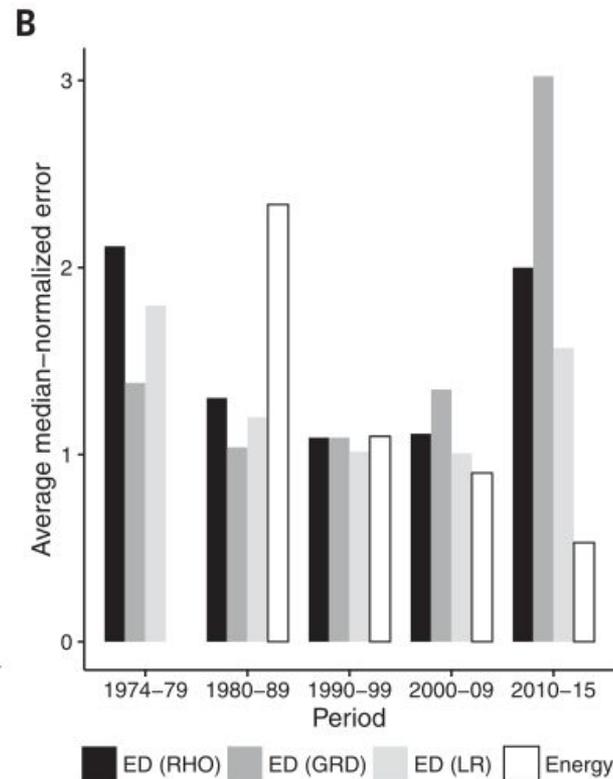
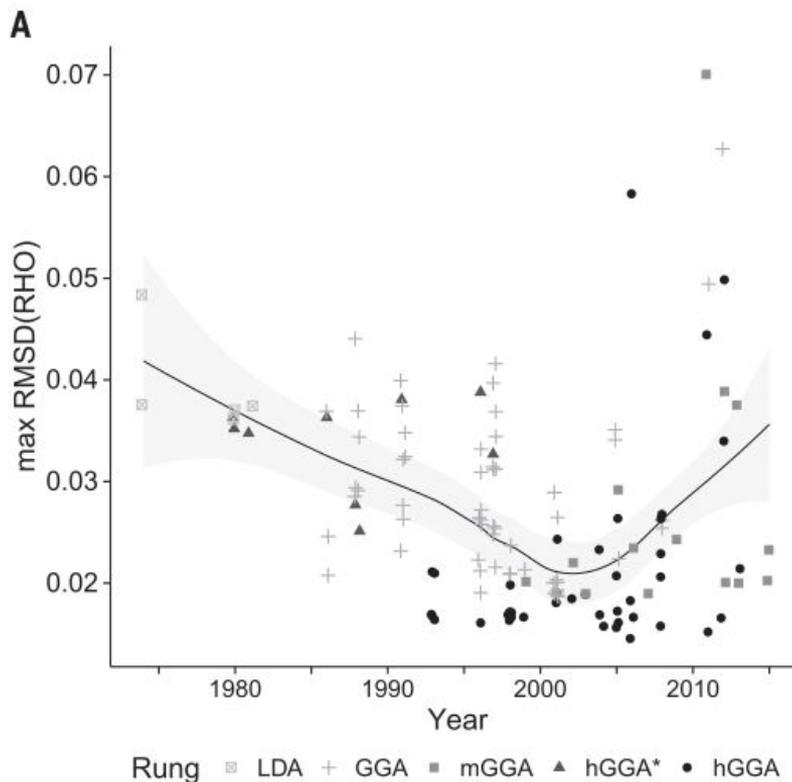


REPORT

THEORETICAL

Density functional theory: the end of the road?

Michael G. Michaelides
John P. Perdew



Measuring the accuracy of the electronic density



Two main groups of methods:

- Metrics based on integrals of the density:

Density moments, Total kinetic energy.

Problems can arise due to possible cancellation of errors!

- Metrics based on point to point measurements. (Local values)

Measuring the accuracy of the electronic density



Some metrics based on point to point comparisons:

$$\text{RMSD} = \sqrt{\frac{1}{N} \sum_i^N \{ \pi r_i^2 [P_{\text{approx}}(r_i) - P_{\text{ref}}(r_i)] \}^2},$$

Root mean-square deviation

$$\text{ISD} = \iiint [P_{\text{approx}}(r) - P_{\text{ref}}(r)]^2 d^3r,$$

Integral squared deviation.

$$\text{IAD} = \iiint |P_{\text{approx}}(r) - P_{\text{ref}}(r)| d^3r.$$

Integral absolute deviation.

P can be any local descriptor such as the density, norm of the gradient, the laplacian, etc...

2. Two theorems for $n(\mathbf{r})$ and $G(\mathbf{r})$ in finite systems.



Kinetic energy density:

$$\tau_{ks} = \frac{1}{2} \sum_i^{\text{occup}} f_i |\nabla \phi_i|^2$$

Kinetic energy per particle:

$$G(\mathbf{r}) = \frac{\tau_{ks}(\mathbf{r})}{n(\mathbf{r})}$$

2. Two theorems for the single particle density in finite systems.



From Kato's theorem (Kato cusp condition): $Z_k = -\frac{a_0}{2n(\mathbf{r})} \frac{dn(\mathbf{r})}{dr} \Big|_{r \rightarrow R_k}$

$$\lim_{r \rightarrow R_k} G(\mathbf{r}) = \frac{Z_k^2}{2}$$

Communications on Pure and Applied Mathematics. **10** (2): 151–177 (1957)

2. Two theorems for the single particle density in finite systems.



From Hoffmann-Ostenhof theorems for the asymptotic behavior of the multi-electronic wave function of finite coulombic systems:

$$\lim_{r \rightarrow \infty} G(\mathbf{r}) = I$$

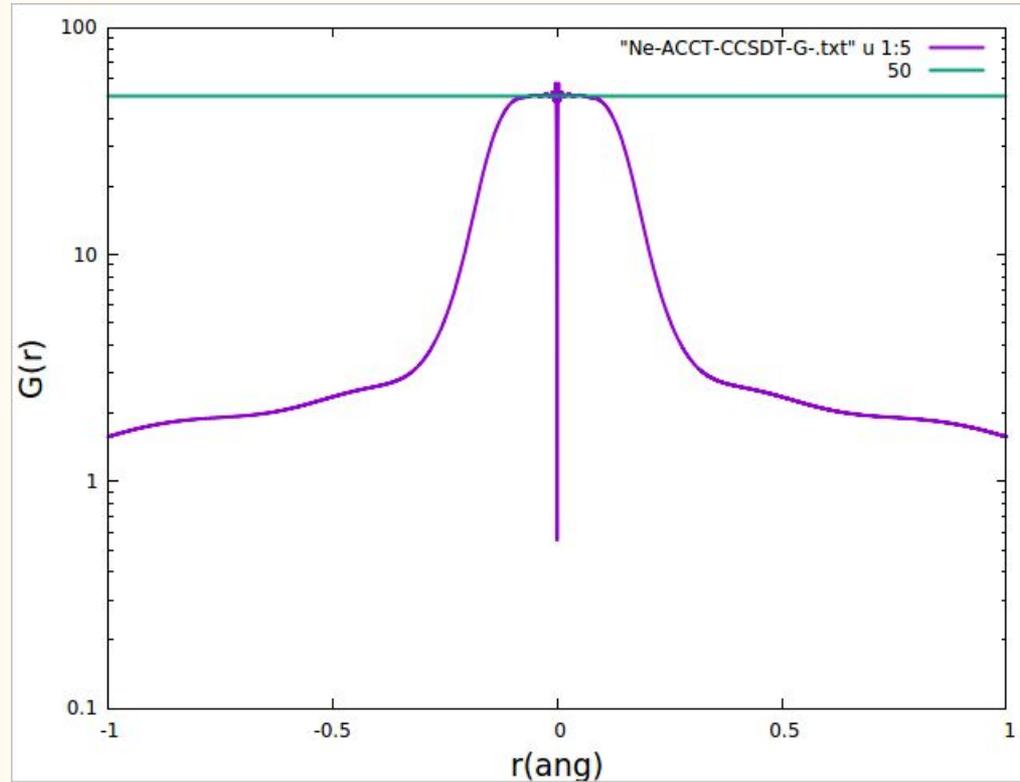
J. Phys.B: Atom. Molec. Phys., **11:17** (1978), *Chem. Phys. Lett.* **258:30** (1996)

3. Computing the asymptotic behavior of $G(\mathbf{r})$ in ABINIT for atoms and molecules.



First some reference calculations: Ne atom CCSDT (ACCT) in GAMESS

$$Z_{\text{Ne}} = 10$$

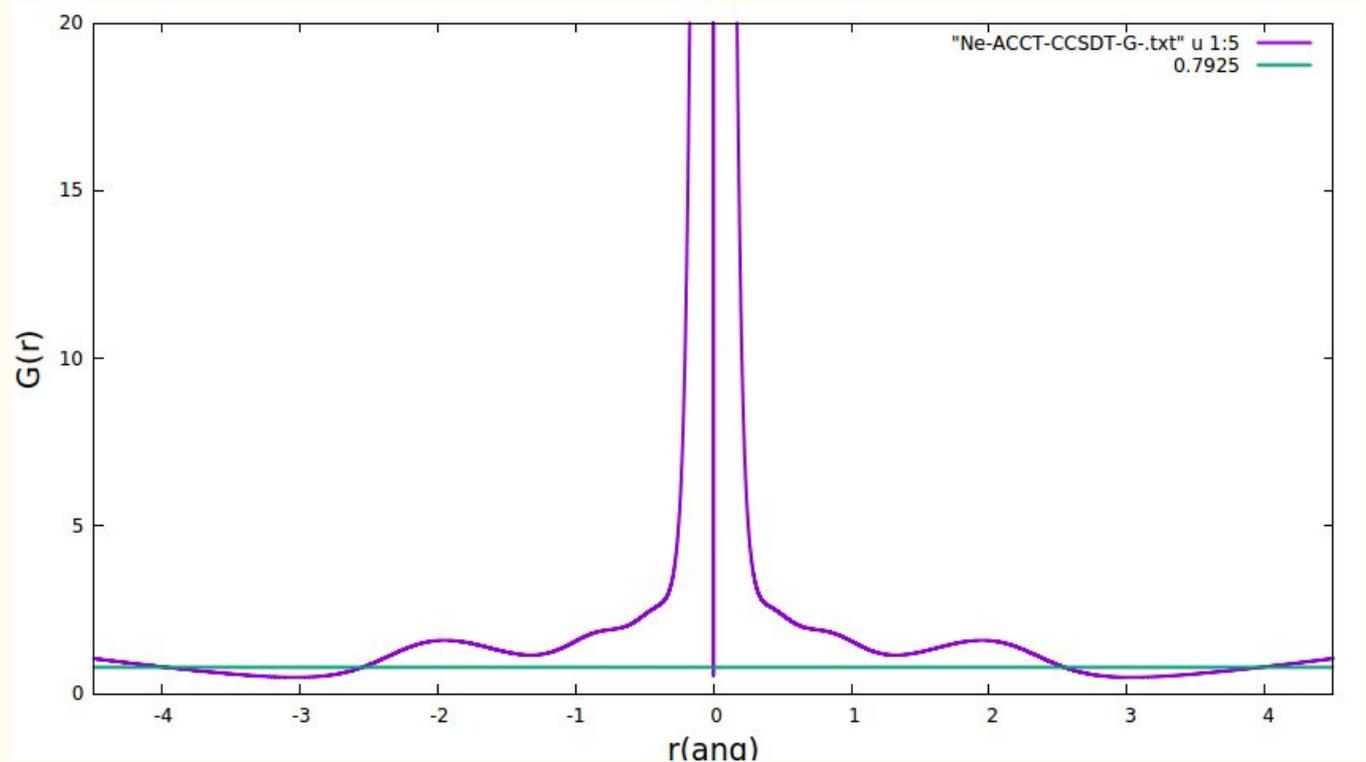


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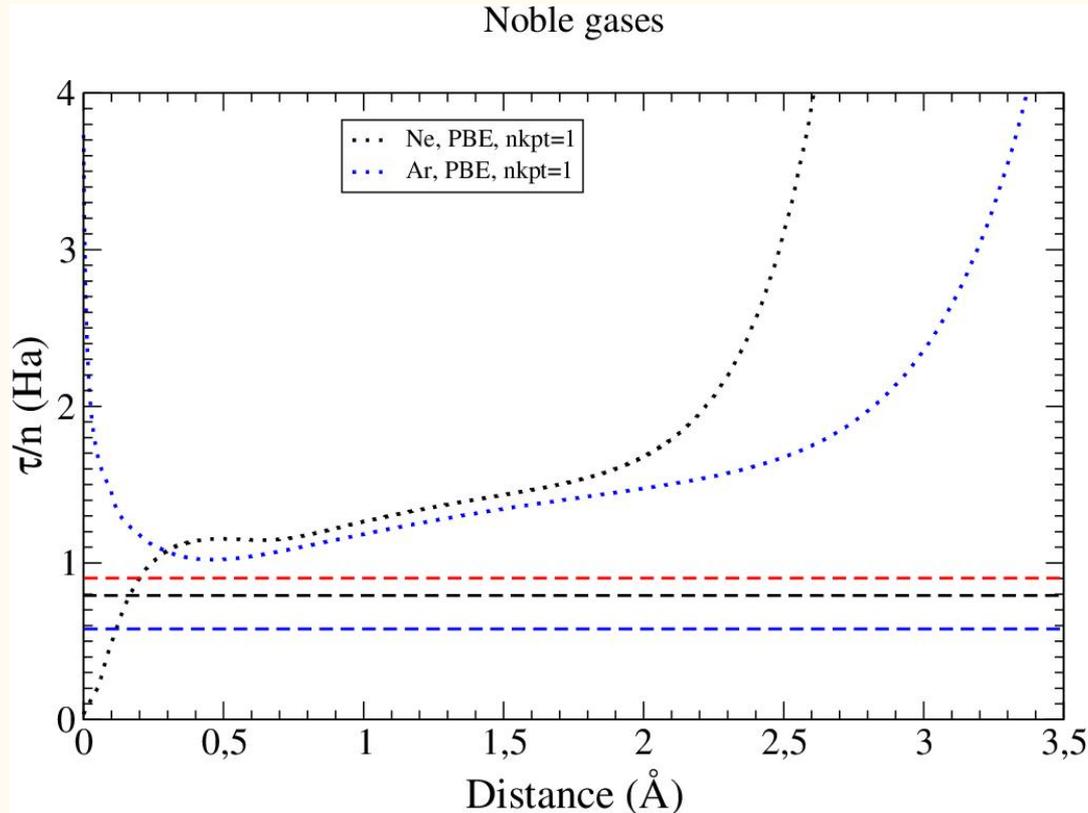
$I=0.7925$ Ha



3. Computing the asymptotic behavior of $G(\mathbf{r})$ in ABINIT for atoms and molecules.



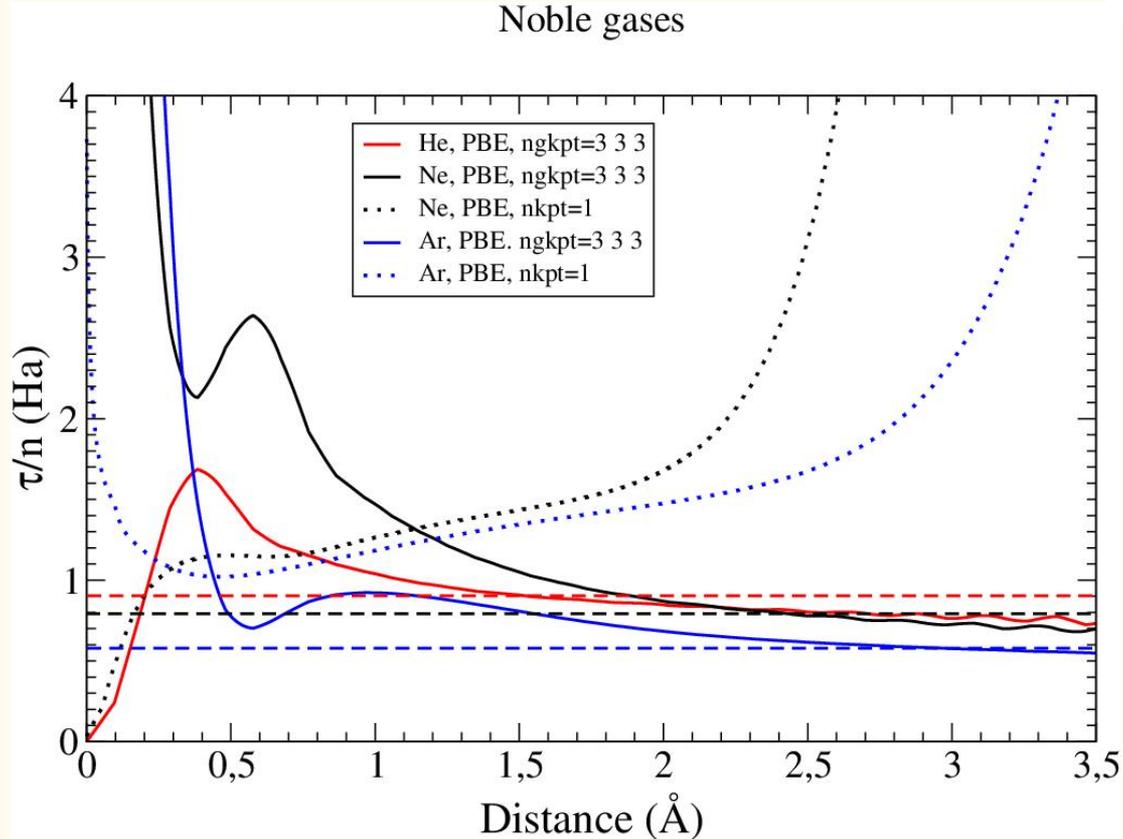
ecut=80 Ha
acell=3*15 A



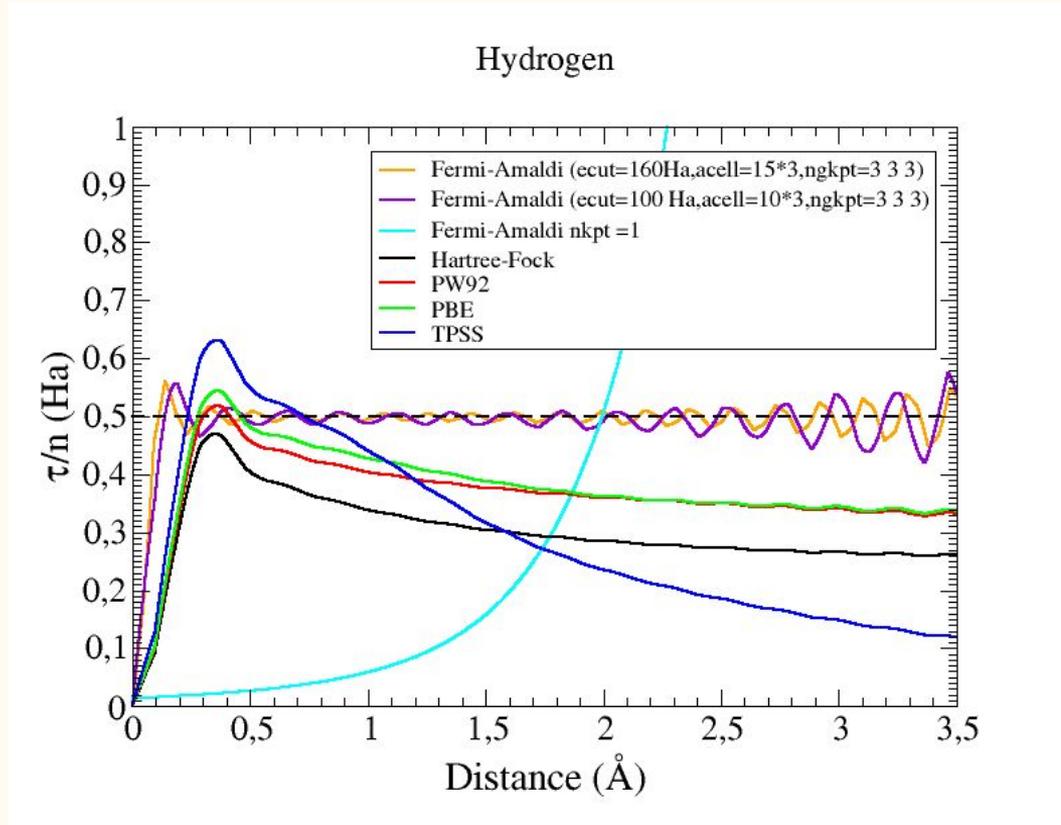
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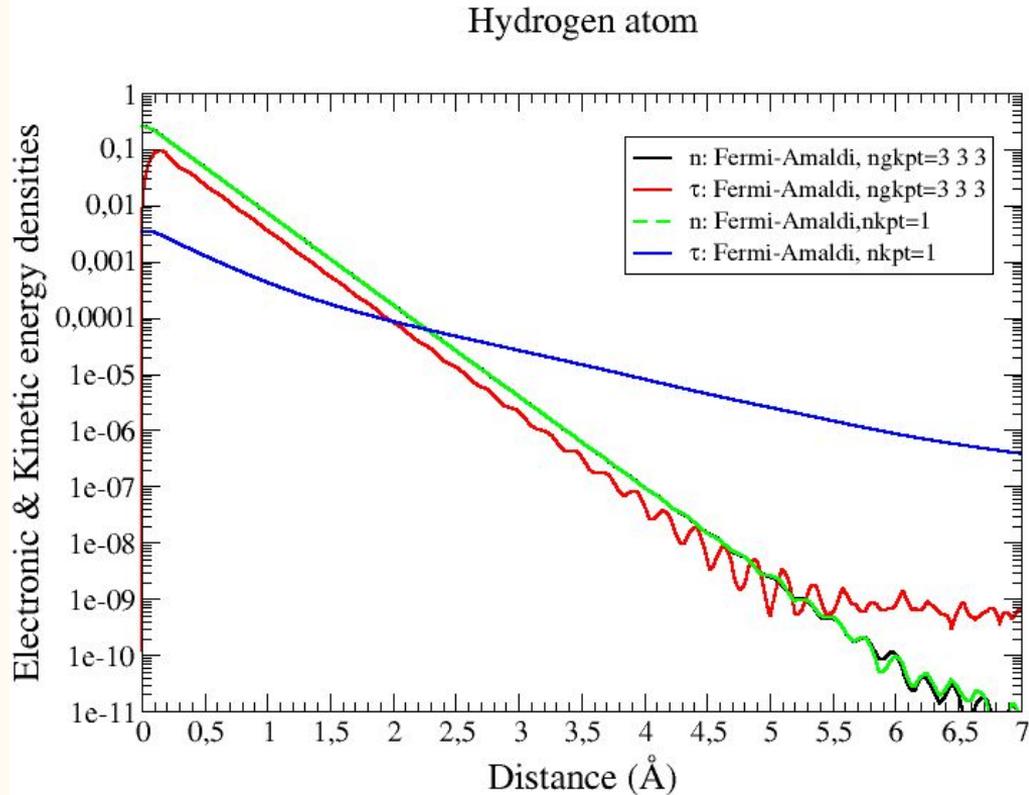
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acell=3*15 A



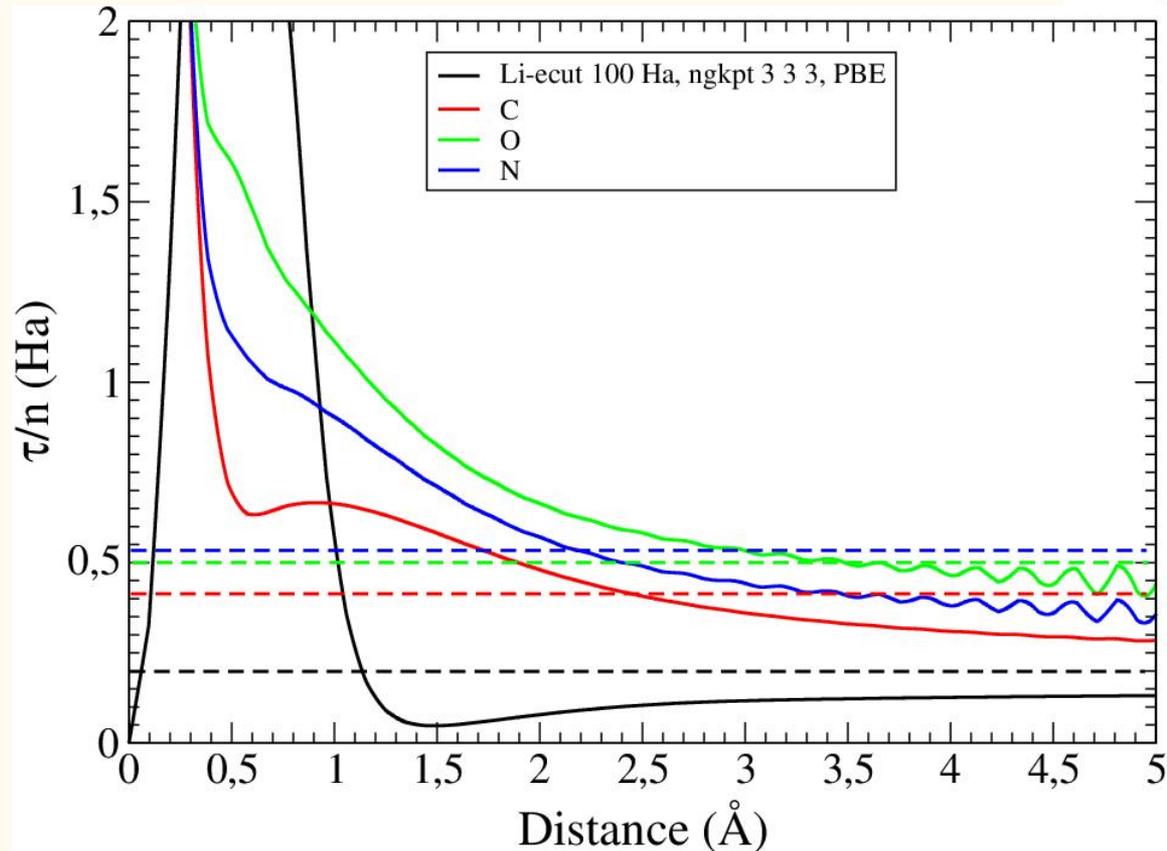
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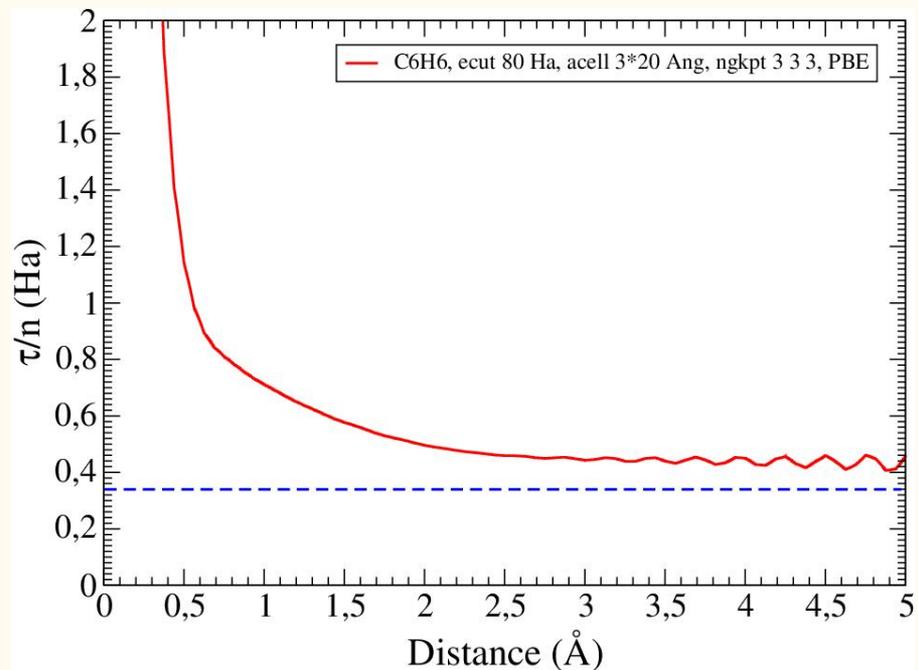
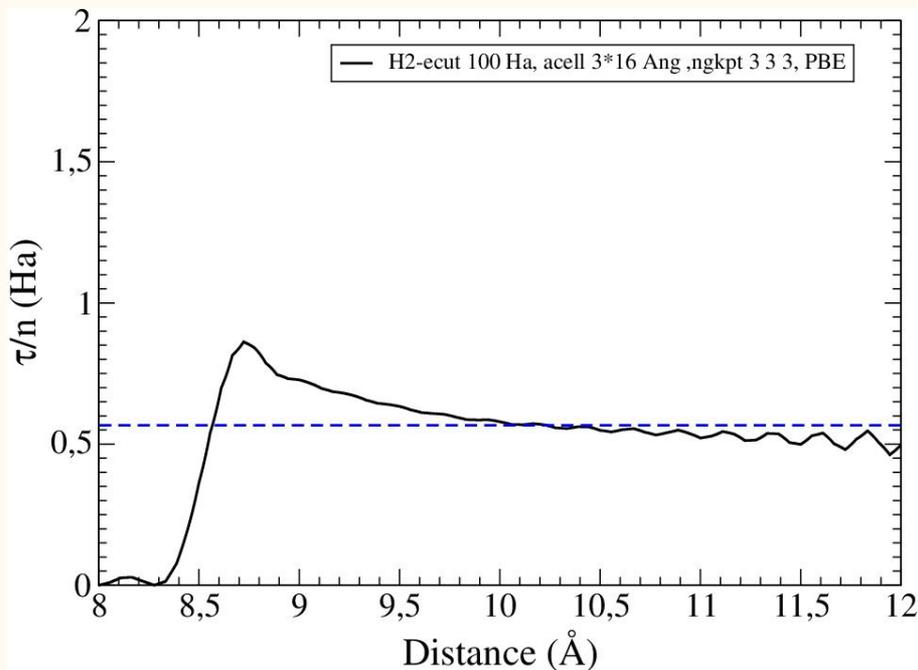
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3. Computing the asymptotic behavior of $G(\mathbf{r})$ in ABINIT for atoms and molecules.



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4. Conclusions and perspectives.

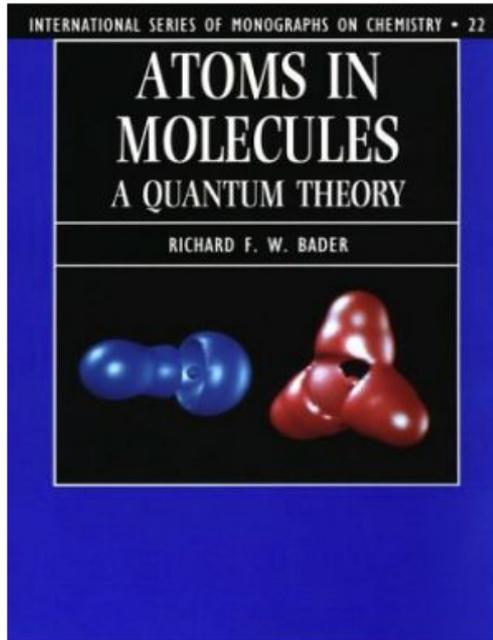
- It is possible to obtain a reasonable behavior of $G(\mathbf{r})$ within the ncpp + plane-waves framework for finite systems if a fine enough k sampling is used.
- Since τ in real space is an input for meta-GGA functionals it is important to add k points convergence studies when dealing with finite systems (i.e. adsorption of molecules on surfaces).
- $G(\mathbf{r})$ can play an important role in the benchmarking of vdW self-consistent functionals from the density point of view.
- What could be the effect of the Coulomb truncation methods (from the last talk) on the asymptotic behavior of G ?
- What is the meaning of the limit of G in extended systems (2D materials, surfaces)??



Thank you for your attention!

Appendix

QTAIM Basics (Bader and Co)



QTAIM Basics (Bader and Co)

$$\nabla\rho = \mathbf{i} \frac{d\rho}{dx} + \mathbf{j} \frac{d\rho}{dy} + \mathbf{k} \frac{d\rho}{dz} \rightarrow \begin{cases} = \vec{\mathbf{0}} & \text{(At critical points and} \\ & \text{at } \infty) \\ \text{Generally } \neq \vec{\mathbf{0}} & \text{(At all other points)} \end{cases}$$

$$\mathbf{A}(\mathbf{r}_c) = \begin{pmatrix} \frac{\partial^2 \rho}{\partial x^2} & \frac{\partial^2 \rho}{\partial x \partial y} & \frac{\partial^2 \rho}{\partial x \partial z} \\ \frac{\partial^2 \rho}{\partial y \partial x} & \frac{\partial^2 \rho}{\partial y^2} & \frac{\partial^2 \rho}{\partial y \partial z} \\ \frac{\partial^2 \rho}{\partial z \partial x} & \frac{\partial^2 \rho}{\partial z \partial y} & \frac{\partial^2 \rho}{\partial z^2} \end{pmatrix}_{\mathbf{r}=\mathbf{r}_c}$$

QTAIM Basics (Bader and Co)

$$\Lambda = \begin{pmatrix} \frac{\partial^2 \rho}{\partial x'^2} & 0 & 0 \\ 0 & \frac{\partial^2 \rho}{\partial y'^2} & 0 \\ 0 & 0 & \frac{\partial^2 \rho}{\partial z'^2} \end{pmatrix}_{\mathbf{r}'=\mathbf{r}_c} = \begin{pmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{pmatrix},$$

$$\nabla^2 \rho(\mathbf{r}) = \nabla \cdot \nabla \rho(\mathbf{r}) = \underbrace{\frac{\partial^2 \rho(\mathbf{r})}{\partial x^2}}_{\lambda_1} + \underbrace{\frac{\partial^2 \rho(\mathbf{r})}{\partial y^2}}_{\lambda_2} + \underbrace{\frac{\partial^2 \rho(\mathbf{r})}{\partial z^2}}_{\lambda_3}$$

QTAIM Basics (Bader and Co)

There are four types of stable critical points having three non-zero eigenvalues:

- $(3, -3)$ Three negative curvatures: ρ is a local maximum.
- $(3, -1)$ Two negative curvatures: ρ is a maximum in the plane defined by the corresponding eigenvectors but is a minimum along the third axis which is perpendicular to this plane.
- $(3, +1)$ Two positive curvatures: ρ is a minimum in the plane defined by the corresponding eigenvectors and a maximum along the third axis which is perpendicular to this plane.
- $(3, +3)$ Three curvatures are positive: ρ is a local minimum.

QTAIM Basics (Bader and Co)

(3,-3)	All curvatures -ve, a local maximum
(3,-1)	Two curvatures are -ve and one is +ve ρ is a maximum in a plane and a minimum perpendicular to this plane - a bond cp
(3,+1)	Two curvatures are +ve and one is -ve ρ is a minimum in a plane and a maximum perpendicular to this plane - a ring cp
3,+3)	All curvatures +ve, a local minimum - a cage cp

QTAIM Basics (Bader and Co)

Poincaré-Hopf relationships:

$$n_{\text{NCP}} - n_{\text{BCP}} + n_{\text{RCP}} - n_{\text{CCP}} = \begin{cases} 1 & \text{(Isolated molecules)} \\ 0 & \text{(Infinite crystals)} \end{cases}$$

QTAIM Basics (Bader and Co)

